**950** Chapter 13: Vector-Valued Functions and Motion in Space

### **Planetary Motion and Satellites 13.6**

In this section, we derive Kepler's laws of planetary motion from Newton's laws of motion and gravitation and discuss the orbits of Earth satellites. The derivation of Kepler's laws from Newton's is one of the triumphs of calculus. It draws on almost everything we have studied so far, including the algebra and geometry of vectors in space, the calculus of vector functions, the solutions of differential equations and initial value problems, and the polar coordinate description of conic sections.



**FIGURE 13.32** The length of **r** is the positive polar coordinate *r* of the point *P*. Thus,  $\mathbf{u}_r$ , which is  $\mathbf{r}/|\mathbf{r}|$ , is also  $\mathbf{r}/r$ . Equations (1) express  $\mathbf{u}_r$  and  $\mathbf{u}_\theta$  in terms of **i** and **j**.

# **Motion in Polar and Cylindrical Coordinates**

When a particle moves along a curve in the polar coordinate plane, we express its position, velocity, and acceleration in terms of the moving unit vectors

$$
\mathbf{u}_r = (\cos \theta)\mathbf{i} + (\sin \theta)\mathbf{j}, \qquad \mathbf{u}_\theta = -(\sin \theta)\mathbf{i} + (\cos \theta)\mathbf{j}, \tag{1}
$$

shown in Figure 13.32. The vector  $\mathbf{u}_r$  points along the position vector  $\overrightarrow{OP}$ , so  $\mathbf{r} = r\mathbf{u}_r$ . The vector  $\mathbf{u}_{\theta}$ , orthogonal to  $\mathbf{u}_r$ , points in the direction of increasing  $\theta$ .

We find from Equations (1) that

$$
\frac{d\mathbf{u}_r}{d\theta} = -(\sin\theta)\mathbf{i} + (\cos\theta)\mathbf{j} = \mathbf{u}_\theta
$$
  

$$
\frac{d\mathbf{u}_\theta}{d\theta} = -(\cos\theta)\mathbf{i} - (\sin\theta)\mathbf{j} = -\mathbf{u}_r.
$$
 (2)

When we differentiate  $\mathbf{u}_r$  and  $\mathbf{u}_\theta$  with respect to *t* to find how they change with time, the Chain Rule gives

$$
\dot{\mathbf{u}}_{r} = \frac{d\mathbf{u}_{r}}{d\theta} \dot{\theta} = \dot{\theta} \mathbf{u}_{\theta}, \qquad \dot{\mathbf{u}}_{\theta} = \frac{d\mathbf{u}_{\theta}}{d\theta} \dot{\theta} = -\dot{\theta} \mathbf{u}_{r}.
$$
 (3)

Hence,

$$
\mathbf{v} = \dot{\mathbf{r}} = \frac{d}{dt} \left( r \mathbf{u}_r \right) = \dot{r} \mathbf{u}_r + r \dot{\mathbf{u}}_r = \dot{r} \mathbf{u}_r + r \dot{\theta} \mathbf{u}_\theta. \tag{4}
$$

See Figure 13.33. As in the previous section, we use Newton's dot notation for time deriva-See Figure 13.33. As in the previous section, we use Newton's dot notation for time derivatives to keep the formulas as simple as we can:  $\dot{u}_r$  means  $d\dot{u}_r/dt$ ,  $\dot{\theta}$  means  $d\theta/dt$ , and so on. The acceleration is #

$$
\mathbf{a} = \dot{\mathbf{v}} = (\ddot{r}\mathbf{u}_r + \dot{r}\dot{\mathbf{u}}_r) + (\dot{r}\dot{\theta}\mathbf{u}_{\theta} + r\ddot{\theta}\mathbf{u}_{\theta} + r\dot{\theta}\dot{\mathbf{u}}_{\theta}).
$$
 (5)

When Equations (3) are used to evaluate  $\dot{\mathbf{u}}_r$  and  $\dot{\mathbf{u}}_\theta$  and the components are separated, the equation for acceleration becomes #

$$
\mathbf{a} = (\ddot{r} - r\dot{\theta}^2)\mathbf{u}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta})\mathbf{u}_\theta.
$$
 (6)

To extend these equations of motion to space, we add zk to the right-hand side of the equation  $\mathbf{r} = r\mathbf{u}_r$ . Then, in these *cylindrical coordinates*,

$$
\mathbf{r} = r\mathbf{u}_r + z\mathbf{k}
$$
  
\n
$$
\mathbf{v} = \dot{r}\mathbf{u}_r + r\dot{\theta}\mathbf{u}_\theta + \dot{z}\mathbf{k}
$$
  
\n
$$
\mathbf{a} = (\ddot{r} - r\dot{\theta}^2)\mathbf{u}_r + (r\dot{\theta} + 2\dot{r}\dot{\theta})\mathbf{u}_\theta + \dot{z}\mathbf{k}.
$$
\n(7)

The vectors  $\mathbf{u}_r$ ,  $\mathbf{u}_\theta$ , and **k** make a right-handed frame (Figure 13.34) in which

$$
\mathbf{u}_r \times \mathbf{u}_\theta = \mathbf{k}, \qquad \mathbf{u}_\theta \times \mathbf{k} = \mathbf{u}_r, \qquad \mathbf{k} \times \mathbf{u}_r = \mathbf{u}_\theta. \tag{8}
$$

### **Planets Move in Planes**

Newton's law of gravitation says that if **r** is the radius vector from the center of a sun of mass  $M$  to the center of a planet of mass  $m$ , then the force **F** of the gravitational attraction between the planet and sun is

$$
\mathbf{F} = -\frac{GmM}{|\mathbf{r}|^2} \frac{\mathbf{r}}{|\mathbf{r}|}
$$
 (9)



**FIGURE 13.33** In polar coordinates, the velocity vector is

 $\mathbf{v} = \dot{r} \mathbf{u}_r + r \dot{\theta} \mathbf{u}_\theta$ # #





**FIGURE 13.34** Position vector and basic unit vectors in cylindrical coordinates.



**FIGURE 13.35** The force of gravity is directed along the line joining the centers of mass.



Combining Equation (9) with Newton's second law,  $\mathbf{F} = m\dot{\mathbf{r}}$ , for the force acting on Combining Equation (9) with Newton's second law,  $\mathbf{F} = m\ddot{\mathbf{r}}$ , for the force acting on the planet gives

$$
m\ddot{\mathbf{r}} = -\frac{GmM}{|\mathbf{r}|^2} \frac{\mathbf{r}}{|\mathbf{r}|},
$$

$$
\ddot{\mathbf{r}} = -\frac{GM}{|\mathbf{r}|^2} \frac{\mathbf{r}}{|\mathbf{r}|}.
$$
(10)

The planet is accelerated toward the sun's center at all times. Equation  $(10)$  says that **r** is a scalar multiple of **r**, so that

> (11)  $\mathbf{r} \times \ddot{\mathbf{r}} = 0.$

A routine calculation shows  $\mathbf{r} \times \mathbf{r}$  to be the derivative of  $\mathbf{r} \times \mathbf{r}$ .

$$
\frac{d}{dt}(\mathbf{r} \times \dot{\mathbf{r}}) = \underbrace{\dot{\mathbf{r}} \times \dot{\mathbf{r}}}_{0} + \mathbf{r} \times \ddot{\mathbf{r}} = \mathbf{r} \times \ddot{\mathbf{r}}.
$$
\n(12)

Hence Equation (11) is equivalent to

$$
\frac{d}{dt}(\mathbf{r} \times \dot{\mathbf{r}}) = \mathbf{0},\tag{13}
$$

which integrates to

$$
\mathbf{r} \times \dot{\mathbf{r}} = \mathbf{C} \tag{14}
$$

for some constant vector **C**.

Equation (14) tells us that **r** and **r** always lie in a plane perpendicular to **C**. Hence, the planet moves in a fixed plane through the center of its sun (Figure 13.36). #

## **Coordinates and Initial Conditions**

We now introduce coordinates in a way that places the origin at the sun's center of mass and makes the plane of the planet's motion the polar coordinate plane. This makes **r** the planet's polar coordinate position vector and makes  $|\mathbf{r}|$  equal to *r* and  $\mathbf{r}/|\mathbf{r}|$  equal to  $\mathbf{u}_r$ . We also position the *z*-axis in a way that makes **k** the direction of **C**. Thus, **k** has the same right-hand relation to  $\mathbf{r} \times \dot{\mathbf{r}}$  that **C** does, and the planet's motion is counterclockwise when right-hand relation to  $\mathbf{r} \times \mathbf{r}$  that C does, and the planet's motion is counterclockwise when viewed from the positive *z*-axis. This makes  $\theta$  increase with *t*, so that  $\dot{\theta} > 0$  for all *t*. Finally, we rotate the polar coordinate plane about the *z*-axis, if necessary, to make the initial ray coincide with the direction **r** has when the planet is closest to the sun. This runs the ray through the planet's **perihelion** position (Figure 13.37). #

If we measure time so that  $t = 0$  at perihelion, we have the following initial conditions for the planet's motion.

- **1.**  $r = r_0$ , the minimum radius, when  $t = 0$
- 2.  $\dot{r} = 0$  when  $t = 0$  (because *r* has a minimum value then)
- **3.**  $\theta = 0$  when  $t = 0$
- **4.**  $|v| = v_0$  when  $t = 0$



**FIGURE 13.36** A planet that obeys Newton's laws of gravitation and motion travels in the plane through the sun's center of mass perpendicular to  $C = r \times \dot{r}$ . #



**FIGURE 13.37** The coordinate system for planetary motion. The motion is counterclockwise when viewed from above, as it is here, and  $\dot{\theta} > 0$ . #

Since

$$
\begin{aligned}\n\mathbf{v}_0 &= |\mathbf{v}|_{t=0} \\
&= |\dot{r} \mathbf{u}_r + r \dot{\theta} \mathbf{u}_\theta|_{t=0} \qquad \text{Equation (4)} \\
&= |r \dot{\theta} \mathbf{u}_\theta|_{t=0} \qquad \dot{r} = 0 \text{ when } t = 0 \\
&= (|r \dot{\theta}| |\mathbf{u}_\theta|)_{t=0} \\
&= |r \dot{\theta}|_{t=0} \qquad |\mathbf{u}_\theta| = 1 \\
&= (r \dot{\theta})_{t=0}, \qquad r \text{ and } \dot{\theta} \text{ both positive}\n\end{aligned}
$$

we also know that

**5.**  $r\dot{\theta} = v_0$  when  $t = 0$ .

## **Kepler's First Law (The Conic Section Law)**

*Kepler's ƒirst law* says that a planet's path is a conic section with the sun at one focus. The eccentricity of the conic is

$$
e = \frac{r_0 v_0^2}{GM} - 1 \tag{15}
$$

and the polar equation is

$$
r = \frac{(1+e)r_0}{1+e\cos\theta}.\tag{16}
$$

The derivation uses Kepler's second law, so we will state and prove the second law before proving the first law.

### **Kepler's Second Law (The Equal Area Law)**

*Kepler's second law* says that the radius vector from the sun to a planet (the vector **r** in our model) sweeps out equal areas in equal times (Figure 13.38). To derive the law, we use Equation (4) to evaluate the cross product  $C = \mathbf{r} \times \dot{\mathbf{r}}$  from Equation (14): !<br>.

$$
\mathbf{C} = \mathbf{r} \times \dot{\mathbf{r}} = \mathbf{r} \times \mathbf{v}
$$
  
=  $r\mathbf{u}_r \times (\dot{r}\mathbf{u}_r + r\dot{\theta}\mathbf{u}_\theta)$  Equation (4)  
=  $r\dot{r}(\mathbf{u}_r \times \mathbf{u}_r) + r(r\dot{\theta})(\mathbf{u}_r \times \mathbf{u}_\theta)$   
=  $r(r\dot{\theta})\mathbf{k}$ . (17)

Setting *t* equal to zero shows that

$$
\mathbf{C} = [r(r\dot{\theta})]_{t=0} \mathbf{k} = r_0 v_0 \mathbf{k}.
$$
 (18)

Substituting this value for **C** in Equation (17) gives #

$$
r_0 v_0 \mathbf{k} = r^2 \dot{\theta} \mathbf{k}, \qquad \text{or} \qquad r^2 \dot{\theta} = r_0 v_0. \tag{19}
$$

This is where the area comes in. The area differential in polar coordinates is

$$
dA = \frac{1}{2}r^2 d\theta
$$



HISTORICAL BIOGRAPHY

Johannes Kepler (1571–1630)

**FIGURE 13.38** The line joining a planet to its sun sweeps over equal areas in equal times.

(Section 10.7). Accordingly,  $dA/dt$  has the constant value

$$
\frac{dA}{dt} = \frac{1}{2}r^2\dot{\theta} = \frac{1}{2}r_0v_0.
$$
 (20)

So  $dA/dt$  is constant, giving Kepler's second law.

For Earth,  $r_0$  is about 150,000,000 km,  $v_0$  is about 30 km/sec, and  $dA/dt$  is about  $2,250,000,000 \text{ km}^2/\text{sec}$ . Every time your heart beats, Earth advances 30 km along its orbit, and the radius joining Earth to the sun sweeps out  $2,250,000,000$  km<sup>2</sup> of area.

### **Proof of Kepler's First Law**

To prove that a planet moves along a conic section with one focus at its sun, we need to express the planet's radius  $r$  as a function of  $\theta$ . This requires a long sequence of calculations and some substitutions that are not altogether obvious.

We begin with the equation that comes from equating the coefficients of  $\mathbf{u}_r = \mathbf{r}/|\mathbf{r}|$  in Equations  $(6)$  and  $(10)$ :

$$
\ddot{r} - r\dot{\theta}^2 = -\frac{GM}{r^2}.\tag{21}
$$

We eliminate  $\dot{\theta}$  temporarily by replacing it with  $r_0v_0/r^2$  from Equation (19) and rearrange the resulting equation to get #

$$
\ddot{r} = \frac{r_0^2 v_0^2}{r^3} - \frac{GM}{r^2}.
$$
 (22)

We change this into a first-order equation by a change of variable. With

$$
p = \frac{dr}{dt}, \qquad \frac{d^2r}{dt^2} = \frac{dp}{dt} = \frac{dp}{dr}\frac{dr}{dt} = p\frac{dp}{dr}, \qquad \text{Chain Rule}
$$

Equation (22) becomes

$$
p\frac{dp}{dr} = \frac{r_0^2 v_0^2}{r^3} - \frac{GM}{r^2}.
$$
 (23)

Multiplying through by 2 and integrating with respect to *r* gives

$$
p^2 = (\dot{r})^2 = -\frac{r_0^2 v_0^2}{r^2} + \frac{2GM}{r} + C_1.
$$
 (24)

The initial conditions that  $r = r_0$  and  $\dot{r} = 0$  when  $t = 0$  determine the value of  $C_1$  to be

$$
C_1 = v_0^2 - \frac{2GM}{r_0}.
$$

Accordingly, Equation (24), after a suitable rearrangement, becomes

$$
\dot{r}^2 = v_0^2 \bigg( 1 - \frac{r_0^2}{r^2} \bigg) + 2GM \bigg( \frac{1}{r} - \frac{1}{r_0} \bigg). \tag{25}
$$

The effect of going from Equation (21) to Equation (25) has been to replace a secondorder differential equation in *r* by a first-order differential equation in *r*. Our goal is still to express *r* in terms of  $\theta$ , so we now bring  $\theta$  back into the picture. To accomplish this, we

divide both sides of Equation (25) by the squares of the corresponding sides of the equadivide both sides of Equation (25) by the squares of the corresponding sides of the equation  $r^2\dot{\theta} = r_0v_0$  (Equation 19) and use the fact that  $\dot{r}/\dot{\theta} = (dr/dt)/(d\theta/dt) = dr/d\theta$  to get !<br>. !<br>.

$$
\frac{1}{r^4} \left(\frac{dr}{d\theta}\right)^2 = \frac{1}{r_0^2} - \frac{1}{r^2} + \frac{2GM}{r_0^2 v_0^2} \left(\frac{1}{r} - \frac{1}{r_0}\right)
$$
\n
$$
= \frac{1}{r_0^2} - \frac{1}{r^2} + 2h\left(\frac{1}{r} - \frac{1}{r_0}\right), \qquad h = \frac{GM}{r_0^2 v_0^2}
$$
\n(26)

To simplify further, we substitute

$$
u = \frac{1}{r}, \qquad u_0 = \frac{1}{r_0}, \qquad \frac{du}{d\theta} = -\frac{1}{r^2}\frac{dr}{d\theta}, \qquad \left(\frac{du}{d\theta}\right)^2 = \frac{1}{r^4}\left(\frac{dr}{d\theta}\right)^2,
$$

obtaining

$$
\left(\frac{du}{d\theta}\right)^2 = u_0^2 - u^2 + 2hu - 2hu_0 = (u_0 - h)^2 - (u - h)^2, \tag{27}
$$

$$
\frac{du}{d\theta} = \pm \sqrt{(u_0 - h)^2 - (u - h)^2}.
$$
 (28)

Which sign do we take? We know that  $\dot{\theta} = r_0 v_0 / r^2$  is positive. Also, *r* starts from a which sign do we take? We know that  $\theta = r_0 v_0/r^2$  is positive. Also, r starts from a<br>minimum value at  $t = 0$ , so it cannot immediately decrease, and  $\dot{r} \ge 0$ , at least for early positive values of *t*. Therefore, #

$$
\frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} \ge 0 \quad \text{and} \quad \frac{du}{d\theta} = -\frac{1}{r^2} \frac{dr}{d\theta} \le 0.
$$

The correct sign for Equation (28) is the negative sign. With this determined, we rearrange Equation (28) and integrate both sides with respect to  $\theta$ :

$$
\frac{-1}{\sqrt{(u_0 - h)^2 - (u - h)^2}} \frac{du}{d\theta} = 1
$$
  

$$
\cos^{-1} \left(\frac{u - h}{u_0 - h}\right) = \theta + C_2.
$$
 (29)

The constant  $C_2$  is zero because  $u = u_0$  when  $\theta = 0$  and  $\cos^{-1}(1) = 0$ . Therefore,

$$
\frac{u-h}{u_0-h}=\cos\theta
$$

and

$$
\frac{1}{r} = u = h + (u_0 - h)\cos\theta.
$$
 (30)

A few more algebraic maneuvers produce the final equation

$$
r = \frac{(1+e)r_0}{1+e\cos\theta},\tag{31}
$$

where

$$
e = \frac{1}{r_0 h} - 1 = \frac{r_0 v_0^2}{GM} - 1.
$$
 (32)

Together, Equations (31) and (32) say that the path of the planet is a conic section with one focus at the sun and with eccentricity  $(r_0v_0^2/GM) - 1$ . This is the modern formulation of Kepler's first law.

# **Kepler's Third Law (The Time–Distance Law)**



The time *T* it takes a planet to go around its sun once is the planet's **orbital period**. *Kepler's third law* says that *T* and the orbit's semimajor axis *a* are related by the equation

$$
\frac{T^2}{a^3} = \frac{4\pi^2}{GM}.
$$
\n(33)

Since the right-hand side of this equation is constant within a given solar system, the ratio of  $T^2$  to  $a^3$  *is the same for every planet in the system.* 

Kepler's third law is the starting point for working out the size of our solar system. It allows the semimajor axis of each planetary orbit to be expressed in astronomical units, Earth's semimajor axis being one unit. The distance between any two planets at any time can then be predicted in astronomical units and all that remains is to find one of these distances in kilometers. This can be done by bouncing radar waves off Venus, for example. The astronomical unit is now known, after a series of such measurements, to be 149,597,870 km.

We derive Kepler's third law by combining two formulas for the area enclosed by the planet's elliptical orbit:

Formula 1: Area = 
$$
\pi ab
$$
  
\nFormula 2: Area =  $\int_0^T dA$   
\n=  $\int_0^T \frac{1}{2} r_0 v_0 dt$  Equation (20)  
\n=  $\frac{1}{2} Tr_0 v_0$ .

Equating these gives

$$
T = \frac{2\pi ab}{r_0 v_0} = \frac{2\pi a^2}{r_0 v_0} \sqrt{1 - e^2}.
$$
 For any ellipse,   

$$
b = a\sqrt{1 - e^2}
$$
 (34)

It remains only to express *a* and *e* in terms of  $r_0$ ,  $v_0$ ,  $G$ , and  $M$ . Equation (32) does this for *e*. For *a*, we observe that setting  $\theta$  equal to  $\pi$  in Equation (31) gives

$$
r_{\max} = r_0 \frac{1+e}{1-e}.
$$

Hence,

$$
2a = r_0 + r_{\text{max}} = \frac{2r_0}{1 - e} = \frac{2r_0 GM}{2GM - r_0 v_0^2}.
$$
 (35)

Squaring both sides of Equation (34) and substituting the results of Equations (32) and (35) now produces Kepler's third law (Exercise 15).



**FIGURE 13.39** The orbit of an Earth satellite:  $2a =$  diameter of Earth + perigee height  $+$  apogee height.

### **Orbit Data**

Although Kepler discovered his laws empirically and stated them only for the six planets known at the time, the modern derivations of Kepler's laws show that they apply to any body driven by a force that obeys an inverse square law like Equation (9). They apply to Halley's comet and the asteroid Icarus. They apply to the moon's orbit about Earth, and they applied to the orbit of the spacecraft *Apollo 8* about the moon.

Tables 13.1 through 13.3 give additional data for planetary orbits and for the orbits of seven of Earth's artificial satellites (Figure 13.39). *Vanguard 1* sent back data that revealed differences between the levels of Earth's oceans and provided the first determination of the precise locations of some of the more isolated Pacific islands. The data also verified that the gravitation of the sun and moon would affect the orbits of Earth's satellites and that solar radiation could exert enough pressure to deform an orbit.



\*Millions of kilometers.







*Syncom 3* is one of a series of U.S. Department of Defense telecommunications satellites. *Tiros II* (for "television infrared observation satellite") is one of a series of weather satellites. *GOES 4* (for "geostationary operational environmental satellite") is one of a series of satellites designed to gather information about Earth's atmosphere. Its orbital period, 1436.2 min, is nearly the same as Earth's rotational period of 1436.1 min, and its orbit is nearly circular  $(e = 0.0003)$ . *Intelsat 5* is a heavy-capacity commercial telecommunications satellite.