NONLINEAR OPTICAL CONSTANTS

H. P. R. Frederikse

The relation between the polarization density *P* of a dielectric medium and the electric field *E* is linear when *E* is small, but becomes nonlinear as *E* acquires values comparable with interatomic electric fields $(10^5 \text{ to } 10^8 \text{ V/cm})$. Under these conditions the relation between *P* and *E* can be expanded in a Taylor's series

$$
P = \varepsilon_0 \chi^{(1)} E + 2 \chi^{(2)} E^2 + 4 \chi^{(3)} E^3 + \cdots
$$
 (1)

where ε _o is the permittivity of free space, while $\chi^{(1)}$ is the linear and $\chi^{(2)}$, $\chi^{(3)}$ etc. the nonlinear optical susceptibilities.

If we consider two optical fields, the first $E_j^{\omega_1}$ (along the *j*-direction at frequency ω_1) and the second $E_k^{\omega_2}$ (along the *k*-direction at frequency ω_2) one can write the second term of the Taylor's series as follows

$$
P_i(\omega_1\omega_2)=2\chi_{ijk}^{\omega_3=\omega_1\pm\omega_2}E_j^{\omega_1}E_k^{\omega_2}
$$

When $\omega_1 \neq \omega_2$ the (parametric) mixing of the two fields gives rise to two new polarizations at the frequencies $\omega_3 = \omega_1 + \omega_2$ and ω_3 ['] = ω₁ – ω₂. When the two frequencies are equal, ω₁ = ω₂ = ω, the result is Second Harmonic Generation (SHG): χ*ijk*(2ω, ω, ω), while equal and opposite frequencies, $\omega_1 = \omega$ and $\omega_2 = -\omega$ leads to Optical Rectification (OR): $\chi_{ijk}(0, \omega, -\omega)$. In the SHG case the following convention is adopted: the second order nonlinear coefficient *d* is equal to one half of the second order nonlinear susceptibility

$$
d_{ijk}=1/2\chi^{(2)}
$$

Because of the symmetry of the indices *j* and *k* one can replace these two by a single index (subscript) *m*. Consequently the notation for the SHG nonlinear coefficient in reduced form is *dim* where *m* takes the values 1 to 6. Only noncentrosymmetric crystals can possess a nonvanishing d_{ik} tensor (third rank). The unit of the SHG coefficients is m/V (in the MKSQ/SI system).

In centrosymmetric media the dominant nonlinearity is of the third order. This effect is represented by the third term in the Taylor's series (Equation 1); it is the result of the interaction of a number of optical fields (one to three) producing a new frequency $ω_4 = ω_1 + ω_2 + ω_3$. The third order polarization is given by

$$
P_j(\omega_1\omega_2\omega_3)=g_4\chi_{jklm}E_k^{\omega_1}E_1^{\omega_2}E_m^{\omega_3}
$$

Third Harmonic Generation (THG) is achieved when $\omega_1 = \omega_2 =$ $ω_3 = ω$. In this case the constant $g_4 = 1/4$. The third order nonlinear coefficient *C* is related to the third order susceptibility as follows:

$$
C_{jklm}=1/4\chi_{jklm}
$$

This coefficient is a fourth rank tensor. In the THG case the matrices must be invariant under permutation of the indices *k, l,* and *m*; as a result the notation for the third order nonlinear coefficient can be simplified to $C_{\!j\!n}^{{}^{}}.$ The unit of $C_{\!j\!n}^{{}^{}}$ is $\mathrm{m}^2\cdot\mathrm{V}^{-2}$ (in the MKSQ/SI system).

Applications of second order nonlinear optical materials include the generation of higher (up to sixth) optical harmonics, the mixing of monochromatic waves to generate sum or difference frequencies (frequency conversion), the use of two monochromatic waves to amplify a third wave (parametric amplification) and the addition of feedback to such an amplifier to create an oscillation (parametric oscillation).

Third order nonlinear optical materials are used for THG, selffocusing, four wave mixing, optical amplification, and optical conjugation. Many of these effects – as well as the variation and modulation of optical propagation caused by mechanical, electric, and magnetic fields (see the preceeding table on "Elasto-Optic, Electro-Optic, and Magneto-Optic Constants") are used in the areas of optical communication, optical computing, and optical imaging.

References

- 1. *Handbook of Laser Science and Technology*, Vol. 111, Part 1; Weber, M. J. Ed., CRC Press, Boca Raton, FL, 1986.
- 2. Dmitriev, V.G., Gurzadyan, G.G., and Nikogosyan, D., *Handbook of Nonlinear Optical Crystals*, Springer-Verlag, Berlin, 1991.
- 3. Shen, Y.R., *The Principles of Nonlinear Optics*, John Wiley, New York, 1984.
- 4. Yariv, A., *Quantum Electronics*, 3rd edition, John Wiley, New York, 1988.
- 5. Bloembergen, N., *Nonlinear Optics,* W.A. Benjamin, New York, 1965.
- 6. Zernike F. and Midwinter, J.E., *Applied Nonlinear Optics*, John Wiley, New York, 1973.
- 7. Hopf, F.A. and Stegeman, G.I., *Applied Classical Electrodynamics*, Volume 2: Nonlinear Optics, John Wiley, New York, 1986.
- 8. *Nonlinear Optical Properties of Organic Molecules and Crystals,* Chemla, D. S., and Zyss, J., Eds., Academic Press, Orlando, FL, 1987.
- 9. *Optical Phase Conjugation*, Fisher, R. A., Ed., Academic Press, New York, 1983.
- 10. Zyss, J., *Molecular Nonlinear Optics: Materials, Devices and Physics*, Academic Press, Boston, 1994.
- 11. Nonlinear Optics, 5 articles in *Physics Today, (Am. Inst. of Phys.)*, Vol. 47, No. 5, May, 1994.

Nonlinear Optical Constants

Selected SHG Coefficients of NLO Crystals*

 $^\ast\,$ These data are taken from References 1 and 2.

Selected THG Coefficients of Some NLO Materials*

 $^\ast\,$ These data are taken from Reference 1.