Stochastic Structural Dynamics

Lecture-7

Random processes-2

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Eergodicty: Temporal and ensemble average

$$
X(t) = a + bt + ct2
$$

\n
$$
\begin{pmatrix} a \\ b \\ c \end{pmatrix} \sim N \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}
$$

\n
$$
\langle X(t) \rangle = \langle a \rangle + \langle b \rangle t + \langle c \rangle t^{2} = 0
$$

\n
$$
\langle X(t_{1}) X(t_{2}) \rangle = a^{2} + b^{2} t_{1} t_{2} + c^{2} t_{1}^{2} t_{2}^{2}
$$

\n
$$
X(t) \sim N(0, a^{2} + b^{2} t^{2} + c^{2} t^{4})
$$

Ergodicty of a random process

Basic notion Equivalence of temporal and ensemble averages

Let $x(t)$ be a sample realization of the random process $X(t)$. We define the time average of a given function of $X(t)$, $g[X(t)]$ by

$$
\mathcal{T}_{\text{av}}\{g[X(t)]\} = \frac{1}{T} \int_0^T g[x(t)] \mathrm{d}t
$$

If $X(t)$ is an ergodic random process, then $g[X(t)] >= \text{T}_{av} \{g[X(t)]\}.$

Definitions

• Ergodicity in mean $X(t)$ is ergodic in mean if

$$
\mathbf{T}_{\text{av}}\{X(t)\}=\frac{1}{T}\int_0^T x(t)\mathrm{d}t=
$$

• Ergodicity in the mean square $X(t)$ is ergodic in meansquare if

$$
T_{av}\lbrace X^2(t)\rbrace = \frac{1}{T}\int_0^T x^2(t)dt =
$$

• Ergodicity in autocorrelation $X(t)$ is said to be ergodic in autocorrelation if

$$
T_{av}\{X(t)X(t1+\tau)\} = \frac{1}{T} \int_0^T x(t)x(t+\tau)dt = \langle X(t)X(t+\tau) \rangle = R_X(\tau)
$$

Remaraks

- 1. The above list of definitions of ergodicity are not exhaustive: several other similar definitions can be constructed by considering other descriptors of the random process.
- 2. Ergodic processes are necessarily stationary in nature; a stationary random process need not be ergodic.
- 3. Physically, ergodicity means that a sufficiently long record of a stationary random process contains all the statistical information about the random phenomenon.

Ergodicity in mean

Let *X(t)* be a stationary random process with specified joint pdf structure

$$
\eta_T = \frac{1}{2T} \int_{-T}^{T} X(t)dt
$$
\n
$$
\Rightarrow \eta_T \text{ is a random variable}
$$
\n
$$
E[\eta_T] = \frac{1}{2T} \int_{-T}^{T} E[X(t)]dt = E[x(t)] = \eta
$$
\n
$$
\sigma_{\eta_T}^2 = \frac{1}{4T^2} \int_{-T-T}^{T-T} E[\{X(t_1) - \eta\} \{X(t_2) - \eta\}]dt_1 dt_2
$$
\n
$$
= \frac{1}{T} \int_{0}^{2T} \left(1 - \frac{\tau}{2T}\right) [R(\tau) - \eta^2] d\tau
$$

Ergodicity in mean

X(t) is said to be ergodic in mean iff

$$
\lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} X(t)dt = E[x(t)] = \eta
$$

$$
\lim_{T \to \infty} \frac{1}{T} \int_{0}^{2T} \left(1 - \frac{\tau}{2T}\right) \left[R(\tau) - \eta^2\right] d\tau \to 0
$$

9

Ergodicity in first order PDF

$$
P_X(x,t) = P_X(x) = P[X(t) \le x]
$$

Define

$$
y(t) = 1 \text{ if } X(t) \le x
$$

$$
y(t) = 0 \text{ if } X(t) > x
$$

$$
\Rightarrow E[y(t)] = 1 \times P[X(t) \le x] + 0 \times P[X(t) > x] = P_X(x)
$$

X(t) is said to be ergodic in first order PDF if *y(t)* is ergodic in mean

Ergodicity in autocorrelation

$$
Define\n\phi(t) = X(t)X(t+\tau)\nE[\phi(t)] = E[X(t)X(t+\tau)] = R_{XX}(\tau)
$$

X(t) is said to be ergodic in autocorrelation if *ϕ(t)* is ergodic in mean

•Criteria for ergodicity in other properties could be developed on similar lines

•The above criteria are applicable if description of the random process is available.

•The notion of ergodicity plays a crucial role in relating observed data to mathematical models of uncertainties

Frequency domain representation of functions of time

Let *x(t)* be a deterministic function of time

Type I : Periodic signals (well behaved) $x(t \pm nT) = x(t)$ $\lim_{t\to\infty}x(t)$ $\lim_{t\to\infty}x(t)$ $\lim_{t\to\infty}x(t)$ Aperiodic signals $\lim |x(t)| \to 0$ Aperiodic signals $\lim |x|$ *t* Aperiodic signals $\lim |x(t)|$ neither goes $\rightarrow \infty$ $\to \infty$ $\to \infty$ **Type II** : Aperiodic signals $\lim |x(t)| \rightarrow$ **Type III** : Aperiodic signals $\lim |x(t)| \to \infty$ **Time signals Type IV :** to zero nor becomes unbounded.

A classification of time signals

Remarks

Realizations of

stationary random

process belong to

Type IV signals.

Type III signals

No hope of any frequency domain representation.

Type I functions

Periodic signals $y(t) = y(t \pm nT)$

Period: the smallest value of T for which the above condition is valid.

$$
y(t) = P\sin \lambda t = P\sin(\lambda t + 2\pi) = P\sin \lambda \left(t + \frac{2\pi}{\lambda}\right) \Rightarrow T = \frac{2\pi}{\lambda}
$$

$$
y(t) = P\cos \lambda t = P\cos(\lambda t + 2\pi) = P\cos \lambda \left(t + \frac{2\pi}{\lambda}\right) \Rightarrow T = \frac{2\pi}{\lambda}
$$

$$
y(t) = P\cos \lambda t + Q\sin \lambda t = P\cos(\lambda t + 2\pi) + Q\sin(\lambda t)
$$

= $P\cos \lambda \left(t + \frac{2\pi}{\lambda}\right) + Q\sin \lambda \left(t + \frac{2\pi}{\lambda}\right) \Rightarrow T = \frac{2\pi}{\lambda}$

$$
y(t) = P\cos 2\lambda t = P\cos(2\lambda t + 2\pi)
$$

$$
= P\cos 2\lambda \left(t + \frac{2\pi}{2\lambda}\right) \Rightarrow T = \frac{\pi}{\lambda}
$$

$$
y(t) = P\cos \lambda t + Q\cos 2\lambda t
$$

= $P\cos(\lambda t + 2\pi) + Q\cos(2\lambda t + 2\pi) \Rightarrow T = \frac{2\pi}{\lambda}$

$$
y(t) = \sum_{n=1}^{N} a_n \cos(\frac{2\pi n}{T}t) + b_n \sin(\frac{2\pi n}{T}t) \Rightarrow
$$

Y(t) is periodic with period=T

According to Fourier's theorem, under general conditions, a periodic function *y(t)* can be represented by

$$
y(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(\frac{2\pi n}{T}t) + b_n \sin(\frac{2\pi n}{T}t)
$$

$$
a_n = \frac{2}{T} \int_{-T/2}^{T/2} y(t) \cos(\frac{2\pi nt}{T}) dt \& b_n = \frac{2}{T} \int_{-T/2}^{T/2} y(t) \sin(\frac{2\pi nt}{T}) dt; n = 1, 2, \dots, \infty
$$

Recall

$$
\cos \theta = \frac{1}{2} [\exp(i\theta) + \exp(-i\theta)] \& \sin \theta = \frac{1}{2i} [\exp(i\theta) - \exp(-i\theta)] \Rightarrow
$$

$$
y(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \left(\frac{\exp\left(i\frac{2\pi nt}{T}\right) + \exp\left(-i\frac{2\pi nt}{T}\right)}{2} \right) + b_n \left(\frac{\exp\left(i\frac{2\pi nt}{T}\right) - \exp\left(-i\frac{2\pi nt}{T}\right)}{2i} \right)
$$

\n
$$
= \frac{a_0}{2} + \frac{1}{2} \sum_{n=1}^{\infty} \exp\left(i\frac{2\pi nt}{T}\right) (a_n - ib_n) + \exp\left(-i\frac{2\pi nt}{T}\right) (a_n + ib_n)
$$

\n
$$
= \sum_{n=-\infty}^{\infty} a_n \exp\left(i\frac{2\pi nt}{T}\right)
$$

\n
$$
\alpha_n = \frac{a_n - ib_n}{2}; \text{ with } a_{-n} = a_n; b_{-n} = -b_n
$$

\n
$$
\alpha_n = \frac{1}{T} \int_{-T/2}^{T/2} y(t) \exp\left(-i\frac{2\pi nt}{T}\right) dt
$$

sine, cosine, amplitude and phase spectra

$$
x(t) \text{ is periodic with period } T
$$

\n
$$
x(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \{a_n \cos \omega_n t + b_n \sin \omega_n t\}; \quad \omega_n = \frac{2\pi n}{T}
$$

\n
$$
a_n = \frac{2}{T} \int_0^T x(t) \cos \omega_n t dt \& \quad b_n = \frac{2}{T} \int_0^T x(t) \sin \omega_n t dt
$$

\n• The plots of a_n and b_n as a function of ω_n are called, respectively, as the Fourier cosine and sine spectra.
\n• The plot of $\sqrt{a_n^2 + b_n^2}$ as a function of ω_n is called the Fourier amplitude spectrum.
\n• The plot of $\tan^{-1} \left(\frac{b_n}{a_n}\right)$ as a function of ω_n is called the Fourier phase spectrum.

Energy and power of a signal

If $x(t)$ is a displacement function, $x^2(t)$ is a Similarly, if $x(t)$ is a velocity function, $x^2(t)$ is a quantity that is proportional to potential energy. quantity that is proportion al to kinetic energy.

 $^{2}\big(t\big)$ $^{2}\big(t\big)$ 0 $\rm 0$ We call $\lim |x^2(t)| dt$ as the total energy in the signal. 1We call $\frac{1}{\pi}$ $\int x^2(t)dt$ as the energy per cycle (power) in the signal. *s s T* $x^2(t)dt$ $x^2(t)dt$ $\frac{1}{T}\int$ $\lim_{\rightarrow\infty}\int$

Total energy and power Discrete power spectrum

Total energy:
$$
\lim_{s \to \infty} \int_{0}^{s} x^2(t) dt \to \infty \Rightarrow
$$

Total energy is not an useful concept.
Energy per cycle= $\frac{1}{T} \int_{0}^{T} x^2(t) dt$ makes sense.
The plot of $\frac{a_n^2 + b_n^2}{2}$ as a function of ω_n is called the discrete power spectrum.
Discrete power spectrum is an useful concept for Type I signals.

Type II signals

$$
x_T(t) = x(t) \text{ for } 0 < t < T
$$

$$
x_T(t + nT) = x(t) \text{ for } n = 1, 2, \dots, \infty
$$

$$
x(t) = \exp(-0.2t)
$$

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$$
x_T(t) = x(t) \text{ for } 0 < t < T
$$

$$
x_T(t + nT) = x(t) \text{ for } n = 1, 2, \dots, \infty
$$

 $\sigma_T(t)$ belongs to Type I of time functions. $x_T(t)$ admits a Fourier series representation. x_{τ} (t $\implies x_r$ (t

Clearly, $\lim_{T \to \infty} x_T(t) \to x(t)$. description of $x_T(t)$ as $T \rightarrow \infty$? Question: What happens to Fourier series based $\lim_{t\to\infty}x_T(t)\to x(t)$ \rightarrow

$$
x_{T}(t) = \sum_{n=-\infty}^{\infty} \alpha_{n} \exp\left(\frac{i2\pi nt}{T}\right)
$$

\n
$$
\alpha_{n} = \frac{1}{T} \int_{-\frac{T}{2}}^{T} x_{T}(t) \exp\left(-\frac{i2\pi nt}{T}\right) dt
$$

\n
$$
x_{T}(t) = \sum_{n=-\infty}^{\infty} \left[\frac{1}{T} \int_{-\frac{T}{2}}^{T} x_{T}(s) \exp\left(-\frac{i2\pi ns}{T}\right) ds \right] \exp\left(\frac{i2\pi nt}{T}\right)
$$

\n
$$
= \sum_{n=-\infty}^{\infty} \left[\frac{1}{T} \int_{-\frac{T}{2}}^{T} x_{T}(s) \exp(-i2\pi nf_{0}s) ds \right] \exp(i2\pi nf_{0}t)
$$

$$
x_{T}(t) = \sum_{n=-\infty}^{\infty} \left[\frac{1}{T} \int_{-T}^{\frac{T}{2}} x_{T}(s) \exp(-i2\pi n f_{0} s) ds \right] \exp(i2\pi n f_{0} t)
$$

\n
$$
f_{0} = \frac{1}{T} \Rightarrow f_{n} = \frac{n}{T} \& \ f_{n+1} = \frac{n+1}{T}
$$

\n
$$
\Rightarrow f_{n+1} - f_{n} = \frac{1}{T} = \Delta f_{n} = \Delta f
$$

\n
$$
x_{T}(t) = \sum_{n=-\infty}^{\infty} \left[\int_{-T}^{\frac{T}{2}} x_{T}(s) \exp(-i2\pi n f_{0} s) ds \right] \exp(i2\pi n f_{0} t) \Delta f_{n}
$$

\n
$$
= \sum_{n=-\infty}^{\infty} X(f_{n}) \exp(i2\pi n f_{0} t) \Delta f_{n}
$$

\n
$$
\lim_{\Delta f_{n} \to \infty} x_{T}(t) \rightarrow x(t) = \int_{-\infty}^{\infty} X(f) \exp(i2\pi f_{n}) df
$$

Definition: Fourier Transform pair

$$
x(t) \text{ is aperiodic; } \lim_{t \to \infty} |x(t)| \to 0
$$

\n
$$
x(t) = \int_{-\infty}^{\infty} X(f) \exp[i2\pi ft] df
$$

\n
$$
X(f) = \int_{-\infty}^{\infty} x(t) \exp[-i2\pi ft] dt
$$

\nPower = $\lim_{T \to \infty} \frac{1}{T} \int_{0}^{T} x^{2}(t) dt \to 0$ and hence not useful.
\nTotal energy = $\lim_{s \to \infty} \int_{0}^{s} x^{2}(t) dt \to \text{ could be an useful quantity.}$
\nEnergy spectrum is an useful concept

$$
x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\omega) \exp(i\omega t) d\omega
$$

$$
X(\omega) = \int_{-\infty}^{\infty} x(t) \exp(-i\omega t) dt
$$

x(t) and X(ω) are said to form a Fourier transform pair

Parseval theorem
\n
$$
\int_{-\infty}^{\infty} x^2(t)dt = \int_{-\infty}^{\infty} x(t) \left[\int_{-\infty}^{\infty} X(f) \exp(i2\pi ft) df \right] dt
$$
\n
$$
= \int_{-\infty}^{\infty} X(f) \left[\int_{-\infty}^{\infty} x(t) \exp(-i2\pi ft) dt \right] df
$$
\n
$$
= \int_{-\infty}^{\infty} X(f) X^*(f) df
$$
\n
$$
\Rightarrow \int_{-\infty}^{\infty} x^2(t) dt = \int_{-\infty}^{\infty} |X(f)|^2 df
$$

Type III time functions
\n
$$
\lim_{t\to\infty} |x(t)| \to \infty
$$

No hope of any frequency domain representations

Type IV

Define
$$
x_T(t) = x(t)
$$
 for $0 < t \leq T$ &
\n
$$
= 0 \text{ for } t > T
$$
\n
$$
X_T(\omega) = \int_{-\infty}^{\infty} x_T(t) \exp(-i\omega t) dt
$$
\n
$$
x_T(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X_T(\omega) \exp(i\omega t) d\omega
$$
\n
$$
\lim_{T \to \infty} \frac{1}{T} \int_{0}^{T} |X_T(f)|^2 df = \text{Total power}
$$
\n
$$
\Rightarrow h(f) = \lim_{T \to \infty} \frac{1}{T} |X_T(f)|^2 = \text{power spectral density function.}
$$
\n
$$
\Rightarrow h(f) df = \text{contribution to the total power made}
$$
\nby the frequency components in the range $(f, f + df)$.

31

Type V: $x(t)$ is a stationary random process

- Let $X(t)$ be a zero mean stationary random process.
- Samples of $X(t)$ belong to Type IV time histories.
- \Rightarrow For each sample the power spectral
- density function can be defined.

: **Definition**

Power spectral density funciton of $X(t)$

$$
S_{XX} (f) = \lim_{T \to \infty} \frac{1}{T} \left\langle \left| X_T (f) \right|^2 \right\rangle
$$

$$
S_{XX}(\omega) = \lim_{T \to \infty} \frac{1}{T} \langle X_T(\omega) X_T^*(\omega) \rangle
$$

\n
$$
= \lim_{T \to \infty} \frac{1}{T} \langle \int_0^T X(t) \exp(-i\omega t) dt \int_0^T X(t) \exp(i\omega t) dt \rangle
$$

\n
$$
= \lim_{T \to \infty} \frac{1}{T} \int_0^T \langle X(t_1) X(t_2) \rangle \exp[i\omega(t_2 - t_1)] dt_1 dt_2
$$

\n
$$
= \lim_{T \to \infty} \frac{1}{T} \int_0^T R_{XX}(t_2 - t_1) \exp[i\omega(t_2 - t_1)] dt_1 dt_2
$$

\n
$$
= \lim_{T \to \infty} \frac{1}{T} \int_{-T}^T [T - |\tau|] R_{XX}(\tau) \exp(i\omega \tau) d\tau
$$

If we restrict our attention to only those $R_{\scriptscriptstyle XX}\left(\tau\right)$ (τ) exp $(i\omega\tau)$ $(\omega) = | R_{XX}(\tau) \exp(i\omega \tau)$ $(\tau) = \frac{1}{2} \int S_{XX}(\omega) \exp(-i\omega \tau)$ $\rm T$ which satisfy the condition 1 \lim_{\longrightarrow} $\int |\tau| R_{XX}(\tau) \exp(i\omega \tau) d\tau \rightarrow 0$, we get the relations exp $\frac{1}{2\pi}\int S_{XX}(\omega)$ exp *T XX T XX XX* XX ^{*(''*) \sim \sim $\frac{1}{2}$ $\frac{1}{2}$} $\frac{1}{T}$ $\int_{0}^{T} |\tau| R_{XX}(\tau) \exp(i\omega \tau) d\tau$ $S_{yy}(\omega) = |R_{yy}(\tau) \exp(i\omega \tau) d\rangle$ $R_{yy}(\tau) = \frac{1}{\tau} \int_{\mathcal{S}} g(x) \exp(-i\omega \tau) d\tau$ τ K_{yy} | τ | exp | $\iota\omega\tau$ | $d\tau$ ω = \pm K_{vv} (τ) exp($\iota \omega \tau$) $d\tau$ τ = $-$ | S_{vv} | ω | exp| $-i\omega\tau$ | $d\omega$ ${\cal T}$ $\rightarrow \infty$ Ξ ∞ $-\infty$ ∞ $-\infty$ \rightarrow Ξ $=$ \Box $_{VV}$ \Box \Box \Box \Box \Box \Box \int \int \int

Remarks
\n(1)
$$
R_{XX}(\tau) = \langle X(t) X(t+\tau) \rangle = \langle X(t) X(t-\tau) \rangle = R_{XX}(-\tau)
$$

\n(2) $S_{XX}(\omega) = \int_{-\infty}^{\infty} R_{XX}(\tau) \exp(i\omega \tau) d\tau$
\n $= \int_{-\infty}^{\infty} R_{XX}(\tau) (\cos \omega \tau + i \sin \omega \tau) d\tau$
\n $= \int_{-\infty}^{\infty} R_{XX}(\tau) \cos \omega \tau d\tau \quad \{\because R_{XX}(\tau) = R_{XX}(-\tau)\}$
\n $= 2 \int_{0}^{\infty} R_{XX}(\tau) \cos \omega \tau d\tau$

Remarks

$$
(3)R_{XX}(0) = \sigma_X^2 = \frac{1}{2\pi} \int_{-\infty}^{\infty} S_{XX}(\omega) d\omega.
$$

 $(4)S_{XX}(\omega)d\omega$ = contribution to the total average power (variance) \Rightarrow Area under the PSD function is the variance of the process. made by frequency components in the range $(\omega,\omega+d\omega)$.

$$
\Rightarrow S_{XX}(\omega) \ge 0
$$

(5) $S_{XX}(-\omega) = \int_{-\infty}^{\infty} R_{XX}(\tau) \exp(-i\omega \tau) d\tau$ (Substitute $s = -\tau$)

$$
= \int_{-\infty}^{\infty} R_{XX}(-s) \exp(i\omega s) ds
$$

$$
= \int_{-\infty}^{\infty} R_{XX}(s) \exp(i\omega s) ds = S_{XX}(\omega)
$$

$(\tau) = \frac{1}{2}$ $S_{XX}(\omega) \exp(-i\omega \tau)$ $\big(\omega \big) \big(\cos\omega \tau + i \sin\omega \tau \big)$ $\big(\mathit{\omega}\big)$ $G_{XX}\left(\omega\right)$ = 2 $S_{XX}\left(\omega\right)$ for ω ≥ 0 01(6) $R_{XX}(\tau) = \frac{1}{2\pi} \int S_{XX}(\omega) \exp$ $=\frac{1}{2\pi}\int S_{XX}(\omega)(\cos \omega \tau + i \sin \theta)$ $=\frac{1}{\pi}\int S_{XX}(\omega)\cos \omega \tau d\omega$ (7) Physical PSD function (defined only for $\omega \ge 0$) $=0$ for $\omega < 0$ $R_{XX}(\tau) = \frac{1}{2} \int S_{XX}(\omega) \exp(-i\omega \tau) d\omega$ $S_{XX}(\omega)(\cos \omega \tau + i \sin \omega \tau) d\omega$ π π π ∞ $-\infty$ ∞ $-\infty$ ∞ $=$ \vert Δ_{vv} \vert ω \vert \in Δ \mathcal{D} \vert $=$ \sum_{VV} ω \int $\cos \omega \tau +$ Ξ \int \int \int **Remarks** Area under $G_{XX}(\omega)$ would still be the variance of

the process.

 $S_{_{XX}}\left(\omega \right)$ 2has properties similar to a pdf $\sigma_{_X}$

PSD

Remarks

(8) Wiener-Khinchine relations
\n
$$
S_{XX}(\omega) = 2 \int_{0}^{\infty} R_{XX}(\tau) \cos \omega \tau d\tau
$$
\n
$$
R_{XX}(\tau) = \frac{1}{\pi} \int_{0}^{\infty} S_{XX}(\omega) \cos \omega \tau d\omega
$$
\n
$$
\Rightarrow
$$
\n
$$
G_{XX}(\omega) = 4 \int_{0}^{\infty} R_{XX}(\tau) \cos \omega \tau d\tau
$$
\n
$$
R_{XX}(\tau) = \frac{1}{2\pi} \int_{0}^{\infty} G_{XX}(\omega) \cos \omega \tau d\omega
$$

A few examples of covariance and psd function pairs

$$
R(\tau) = \frac{1}{2\pi} \int_{-\infty}^{\infty} S(\omega) \exp(j\omega\tau) d\omega \quad S(\omega) = \int_{-\infty}^{\infty} R(\tau) \exp(-j\omega\tau) d\omega
$$

\n
$$
\frac{\partial(\tau)}{\partial(\tau)} = \frac{1}{2\pi\delta(\omega - \beta)}
$$

\n
$$
\frac{\partial(\omega - \beta)}{\partial(\omega - \beta)}
$$

\n
$$
\exp(-\alpha|\tau|)
$$

\n
$$
\exp(-\alpha|\tau|)
$$

\n
$$
\exp(-\alpha\tau^2)
$$

\n
$$
\exp(-\alpha|\tau|) \cos \beta\tau
$$

\n
$$
\frac{\alpha}{\alpha^2 + (\omega - \beta)^2} + \frac{\alpha}{\alpha^2 + (\omega + \beta)^2}
$$

\n
$$
\frac{2\alpha}{\alpha^2 + (\omega - \beta)^2} + \frac{\alpha}{\alpha^2 + (\omega + \beta)^2}
$$

\n
$$
2\exp(-\alpha\tau^2) \cos \beta\tau \quad \sqrt{\frac{\pi}{\alpha}} \left\{ \exp\left(-\frac{(\omega - \beta)^2}{4\alpha}\right) + \exp\left(-\frac{(\omega + \beta)^2}{4\alpha}\right) + \right\}
$$

\n
$$
\frac{\sin \sigma\tau}{\pi\tau}
$$

\n
$$
\frac{\int_{0}^{1} |\omega| < \sigma}{|\omega| > \sigma}
$$

Typical psd function of wind velocity

Typical psd function of waves

$$
G(\omega) = c_0 \frac{1}{\omega^5} \exp\left(-\frac{c_1}{\omega^4}\right)
$$

Typical psd function of earthquake ground acceleration

Evolutionary random process

 $X(t) = V_1(t)$ if $0 < t < t_1$ $\& X(t) = V_2(t) \text{ if } t > t_1.$ Let $V_1(t)$ & $V_2(t)$ be zero mean, stationary random processes $\left\{ S_{VVi}\big(\omega\big)\right\} _{i=1}^{2}$ with psd functions $\left\{ S_{VVi}(\omega) \right\}_{i=1}^{r}$. Consider a random process $X(t)$ defined as \Rightarrow We can write S_{VU} (ω \equiv $S_{XX}(a,t) = S_{VV_1}(a)$ if $0 < t < t_1$ $S_{XX}(\omega, t) = S_{VV_2}(\omega) \text{ if } t > t_1$ This notion can be generalized to define nonstationary random processes with time dependent psd functions. Such processes are called as evolutionary random proc esses.