

Stochastic Structural Dynamics

Lecture-24

Markov Vector Approach-4

Dr C S Manohar

Department of Civil Engineering

Professor of Structural Engineering

Indian Institute of Science

Bangalore 560 012 India

manohar@civil.iisc.ernet.in



Solution of
FPK equations

$$M\ddot{X} + C\dot{X} + KX = \Gamma W(t); t \geq 0; X(0) = X_0; \dot{X}(0) = \dot{X}_0$$
$$X(t) \sim N \times 1; \Gamma \sim n \times m; W(t) \sim m \times 1$$
$$\langle W(t) \rangle = 0; \langle W(t)W^t(t+\tau) \rangle = [2D_{ij}]_{m \times m} \delta(\tau)$$

Transient & Steady
state solutions

$$\dot{x} + \beta(x) = w(t); t \geq 0 \text{ \& } x(0) = x_0$$
$$\langle w(t) \rangle = 0; \langle w(t_1)w(t_2) \rangle = 2D\delta(t_1 - t_2)$$

$$\ddot{x} + \dot{x}f(H) + g(x) = w(t); t \geq 0; x(0) = x_0; \dot{x}(0) = \dot{x}_0.$$
$$\langle w(t) \rangle = 0; \langle w(t)w(t+\tau) \rangle = 2D\delta(\tau)$$

Steady state
solutions

$$m_j \ddot{X}_j + m_j \beta_j \dot{X}_j + \frac{\partial U}{\partial X_j} = W_j(t); t \geq 0; X_j(0) = x_{j0} \text{ \& } \dot{X}_j(0) = \dot{x}_{j0}$$
$$U = \frac{1}{2} X^t KX; \langle W_j(t) \rangle = 0; \langle W_j(t)W_k(t+\tau) \rangle = 2D_{jk} \delta_{jk} \delta(\tau);$$
$$\frac{m_j \beta_j}{2D_{jj}} = \gamma \forall j = 1, 2, \dots, n$$

Moment equations

$$dX(t) = f[t, X(t)]dt + G[t, X(t)]dB(t); t \geq 0; X(0) = X_0$$

$$\frac{\partial p}{\partial t} = -\sum_{j=1}^n \frac{\partial}{\partial x_j} [f_j(x, t) p] + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2}{\partial x_i \partial x_j} [(GDG^t)_{ij} p]$$

$$\frac{d}{dt} \langle h[X(t), t] \rangle = \left\langle \frac{\partial h}{\partial t} \right\rangle + \sum_{j=1}^n \left\langle f_j(X, t) \frac{\partial h}{\partial X_j} \right\rangle + \sum_{i=1}^n \sum_{j=1}^n \left\langle (GDG^t)_{ij} \frac{\partial^2 h}{\partial X_i \partial X_j} \right\rangle$$

$$\dot{X} + \beta X = w(t); t \geq 0; X(0) = X_0$$

$$\ddot{X} + 2\eta\omega\dot{X} + \omega^2 X = w(t); t \geq 0; X(0) \& \dot{X}(0) \text{ specified}$$

Example

$$\dot{X} + \beta X + \alpha X^3 = w(t); t \geq 0; X(0) = X_0$$

$$dX(t) = -\beta X dt - \alpha X^3 dt + dB(t)$$

$$M_k = \langle X^k(t) \rangle$$

$$f = -\beta X - \alpha X^3$$

$$\frac{d}{dt} \langle h[X(t), t] \rangle = \left\langle \frac{\partial h}{\partial t} \right\rangle + \left\langle -(\beta X + \alpha X^3) \frac{\partial h}{\partial X_j} \right\rangle + D \left\langle \frac{\partial^2 h}{\partial X^2} \right\rangle$$

$$\dot{m}_1 = \langle [-\beta X - \alpha X^3] \rangle = -\beta m_1 - \alpha m_3$$

$$\dot{m}_2 = \langle [-\beta X - \alpha X^3] 2X \rangle + 2D = -2\beta m_2 - 2\alpha m_4 + 2D$$

$$\dot{m}_3 = \langle [-\beta X - \alpha X^3] 3X^2 \rangle + 6D m_1 = -3\beta m_3 - 3\alpha m_5 + 6D m_1$$

$$\dot{m}_4 = \langle [-\beta X - \alpha X^3] 4X^3 \rangle + 12D m_2 = -4\beta m_4 - 4\alpha m_6 + 12D m_2$$

⋮

Remarks

- Moment equations form an infinite hierarchy of equations which at no stage provide sufficient number of equations to solve for the moments: **CLOSURE PROBLEM**
- This is the characteristic feature of moment equations of non linear systems driven by random excitations
- Closure approximations
 - Assume that the higher order moments, beyond a given order, are related to the lower order ones as though they obey an adhocly specified pdf (Ex: Gaussian closure)
 - Neglect cumulants beyond a specified order

Note : This problem occurs for non - white excitation also

$$m\ddot{x} + c\dot{x} + kx + \alpha x^3 = f(t); t \geq 0; x(0) \text{ \& } \dot{x}(0) \text{ specified}$$

Taking expectations on both sides

$$m\langle \ddot{x}(t) \rangle + c\langle \dot{x}(t) \rangle + k\langle x(t) \rangle + \alpha\langle x^3(t) \rangle = \langle f(t) \rangle$$

Equation for mean response contains 3rd order expectations.

$$m\langle x(t_1)\ddot{x}(t) \rangle + c\langle x(t_1)\dot{x}(t) \rangle + k\langle x(t_1)x(t) \rangle + \alpha\langle x(t_1)x^3(t) \rangle = \langle x(t_1)f(t) \rangle$$

$$m\frac{\partial^2 R_{xx}(t, t_1)}{\partial t^2} + c\frac{\partial R_{xx}(t, t_1)}{\partial t} + kR_{xx}(t, t_1) + \alpha\langle x(t_1)x^3(t) \rangle = R_{xf}(t_1, t)$$

Equation for ACF of $x(t)$ contains $\alpha\langle x(t_1)x^3(t) \rangle$

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Closure problem

Example

$$\ddot{X} + a_2 \dot{X} + [a_1 + w_1(t)] X = w_2(t);$$

$t \geq 0$; $X(0)$ & $\dot{X}(0)$ specified.

$$dX_1 = X_2 dt$$

$$dX_2 = -a_2 X_2 dt - a_1 X_1 dt - X_1 dB_1(t) + dB_2(t)$$

$$\begin{Bmatrix} dX_1 \\ dX_2 \end{Bmatrix} = \begin{Bmatrix} X_2 \\ -a_2 X_2 - a_1 X_1 \end{Bmatrix} dt + \begin{bmatrix} 0 & 0 \\ -X_1 & 1 \end{bmatrix} \begin{Bmatrix} dB_1(t) \\ dB_2(t) \end{Bmatrix}$$

$$\langle dB_i(t) dB_j(t + \tau) \rangle = 2D_{ij} \delta_{ij} \delta(\tau); i, j = 1, 2$$

$$GDG^t = \begin{bmatrix} 0 & 0 \\ -X_1 & 1 \end{bmatrix} \begin{bmatrix} D_{11} & 0 \\ 0 & D_{22} \end{bmatrix} \begin{bmatrix} 0 & 0 \\ -X_1 & 1 \end{bmatrix}^t = \begin{bmatrix} 0 & 0 \\ 0 & D_{11} X_1^2 + D_{22} \end{bmatrix}$$

$$\begin{Bmatrix} dX_1 \\ dX_2 \end{Bmatrix} = \begin{Bmatrix} X_2 \\ -a_2 X_2 - a_1 X_1 \end{Bmatrix} dt + \begin{bmatrix} 0 & 0 \\ -X_1 & 1 \end{bmatrix} \begin{Bmatrix} dB_1(t) \\ dB_2(t) \end{Bmatrix}$$

$$M_{g,IS} = \langle \overset{g}{X(t)} \overset{f}{\dot{X}(t)} \rangle$$

$$\dot{m}_{10} = \langle X_2 \rangle = m_{01}$$

$$\dot{m}_{01} = \langle [-a_2 X_2 - a_1 X_1] \rangle = -a_2 m_{01} - a_1 m_{10}$$

$$\dot{m}_{20} = \langle X_2 2X_1 \rangle = 2m_{11}$$

$$\dot{m}_{11} = \langle X_2 X_2 + [-a_2 X_2 - a_1 X_1] X_1 \rangle = m_{02} - a_2 m_{11} - a_1 m_{20}$$

$$\dot{m}_{20} = \langle [-a_2 X_2 - a_1 X_1] 2X_2 \rangle + \langle X_1^2 D_{11} + D_{22} \rangle$$

$$= -2a_2 m_{02} - 2a_1 m_{11} + D_{11} m_{20} + D_{22}$$

I order moments

$$\dot{m}_{10} = \langle X_2 \rangle = m_{01}$$

$$\dot{m}_{01} = \langle [-a_2 X_2 - a_1 X_1] \rangle = -a_2 m_{01} - a_1 m_{10}$$

II order moments

$$\dot{m}_{20} = \langle X_2 2X_1 \rangle = \underline{2m_{11}}$$

$$\dot{m}_{11} = \langle X_2 X_2 + [-a_2 X_2 - a_1 X_1] X_1 \rangle = \underline{m_{02}} - a_2 \underline{m_{11}} - a_1 m_{20}$$

$$\begin{aligned} \dot{m}_{20} &= \langle [-a_2 X_2 - a_1 X_1] 2X_2 \rangle + \langle X_1^2 D_{11} + D_{22} \rangle \\ &= -2a_2 \underline{m_{02}} - 2a_1 \underline{m_{11}} + D_{11} \underline{m_{20}} + D_{22} \end{aligned}$$

Moment equations are closed & hence can be evaluated exactly (numerically at least)

Steady state

$$\dot{m}_{10} = 0 \Rightarrow m_{01} = 0$$

$$\dot{m}_{01} = 0 \Rightarrow -a_2 m_{01} - a_1 m_{10} = 0$$

$$m_{01} = 0 \ \& \ m_{10} = 0$$

$$\dot{m}_{20} = 0 \Rightarrow 2m_{11} = 0 \Rightarrow m_{11} = 0$$

$$\dot{m}_{11} = 0 \Rightarrow m_{02} - a_2 m_{11} - a_1 m_{20} = 0 \Rightarrow m_{02} - a_1 m_{20} = 0$$

$$\dot{m}_{20} = 0 \Rightarrow -2a_2 m_{02} - 2a_1 m_{11} + D_{11} m_{20} + D_{22} = 0$$

$$-2a_2 m_{02} + D_{11} m_{20} + D_{22} = 0$$

⋮

Question: Are the steady state moments realizable?

That is. are these moments stable?

Strategy : perturb the solutions and see if perturbations die out.

The moment equations can also be used to study stability of the system in terms of response moments.

$$\dot{Y} = AY + B$$

$$\tilde{Y} = Y + \nu$$

$\nu(t)$ = small perturbation

\Rightarrow

$$\dot{Y} + \dot{\nu} = A(Y + \nu) + B$$

$$\Rightarrow \dot{\nu} = A\nu \Rightarrow \nu(t) = \nu_0 \exp(st) \Rightarrow A\nu_0 = sI\nu_0$$

$$|A - sI| = 0$$

Perturbations do not grow in time if the real part of the eigenvalues of A are all ≤ 0 .

Equations for two time moments

Consider $n = 2$

$$dX(t) = f[t, X(t)]dt + G[t, X(t)]dB(t); t \geq 0; X(0) = X_0$$

$$\frac{\partial p}{\partial t} = -\sum_{j=1}^n \frac{\partial}{\partial x_j} [f_j(x, t) p] + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2}{\partial x_i \partial x_j} [(GDG^t)_{ij} p]$$

$$p = p(\tilde{x}; t | \tilde{x}_0; t_0) = p(x_1, x_2; t | x_{01}, x_{02}; t_0)$$

Consider

$$\langle X_1^m(t) X_1^n(t_0) \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \underline{x_1^m \eta_1^n p(x_1, \eta_1; t, t_0)} dx_1 d\eta_1$$

$$\underline{p(x_1, \eta_1; t, t_0)} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \underline{p(x_1, x_2, \eta_1, \eta_2; t, t, t_0, t_0)} dx_1 dx_2 d\eta_1 d\eta_2$$

$$\langle X_1^m(t) X_1^n(t_0) \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_1^m \eta_1^n p(x_1, \eta_1; t, t_0) dx_1 d\eta_1$$

$$p(x_1, \eta_1; t, t_0) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} p(x_1, x_2, \eta_1, \eta_2; t, t, t_0, t_0) dx_2 d\eta_2$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} p(x_1, x_2; t, t_0 | \eta_1, \eta_2; t, t_0) p(\eta_1, \eta_2; t, t_0) dx_2 d\eta_2$$

$$\langle X_1^m(t) X_1^n(t_0) \rangle =$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_1^m \eta_1^n p(x_1, x_2; t, t_0 | \eta_1, \eta_2; t, t_0) p(\eta_1, \eta_2; t, t_0) dx_1 d\eta_1 dx_2 d\eta_2$$

$$v(x_1, x_2; t, t_0) =$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \eta_1^n p(x_1, x_2; t, t_0 | \eta_1, \eta_2; t, t_0) p(\eta_1, \eta_2; t, t_0) d\eta_1 d\eta_2$$

$$\langle X_1^m(t) X_1^m(t_0) \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x_1^m \underline{v(x_1, x_2; t, t_0)} dx_1 dx_2$$

$$\underline{v(x_1, x_2; t, t_0)} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \eta_1^n p(x_1, x_2; t, t_0 | \eta_1, \eta_2; t, t_0) p(\eta_1, \eta_2; t, t_0) d\eta_1 d\eta_2$$

Consider the FPK equation

$$\frac{\partial p}{\partial t} = - \sum_{j=1}^n \frac{\partial}{\partial x_j} [f_j(x, t) p] + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2}{\partial x_i \partial x_j} [(GDG^t)_{ij} p]$$

It follows that

$$\frac{\partial v}{\partial t} = - \sum_{j=1}^n \frac{\partial}{\partial x_j} [f_j(x, t) v] + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2}{\partial x_i \partial x_j} [(GDG^t)_{ij} v]$$

$$\lim_{t \rightarrow t_0} v(x_1, x_2; t, t_0) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \eta_1^n \delta(x_1 - \eta_1) \delta(x_2 - \eta_2) p(\eta_1, \eta_2; t_0, t_0) d\eta_1 d\eta_2$$

$$= x_1^n p(y_1, y_2; t_0, t_0)$$

Backward Kolmogorov equation and reliability function

References :

- Y K Lin and G C Cai, 1995, Probabilistic structural dynamics, McGraw Hill, NY]
- D R Cox and H D Miller, 1965, The theory of stochastic processes, Chapman and Hall, NY.

Backward Kolmogorov equation and reliability function

Let $X(t)$ be a scalar Markov process. For $t_1 < t_2 < t_3$ we have

$$p(x_3; t_3 | x_1; t_1) = \int_{-\infty}^{\infty} \underbrace{p(x_3; t_3 | x_2; t_2)} p(x_2; t_2 | x_1; t_1) dx_2$$

$$p(x_3; t_3 | \underline{x_2}; t_2) = p(x_3; t_3 | \underline{x_1} + x_2 - \underline{x_1}; t_2)$$

$$= p(x_3; t_3 | x_1; t_2) + (x_2 - x_1) \left. \frac{\partial p}{\partial x_2} \right|_{x_2=x_1}$$

$$+ \frac{1}{2} (x_2 - x_1)^2 \left. \frac{\partial^2 p}{\partial x_2^2} \right|_{x_2=x_1} + o\left[(x_2 - x_1)^2\right]$$

$$\begin{aligned}
p(x_3; t_3 | x_1; t_1) &= \int_{-\infty}^{\infty} p(x_3; t_3 | x_2; t_2) p(x_2; t_2 | x_1; t_1) dx_2 \\
&= p(x_3; t_3 | x_1; t_2) + \left. \frac{\partial p}{\partial x_2} \right|_{x_2=x_1} \int_{-\infty}^{\infty} (x_2 - x_1) p(x_2; t_2 | x_1; t_1) dx_2 \\
&\quad + \left. \frac{1}{2} \frac{\partial^2 p}{\partial x_2^2} \right|_{x_2=x_1} \int_{-\infty}^{\infty} (x_2 - x_1)^2 p(x_2; t_2 | x_1; t_1) dx_2 + o\left[(x_2 - x_1)^2\right] \\
&\Rightarrow \frac{1}{\Delta t} \left[p(x_3; t_3 | x_1; t_2) - p(x_3; t_3 | x_1; t_1) \right] \\
&\quad + \frac{1}{\Delta t} \left. \frac{\partial p}{\partial x_2} \right|_{x_2=x_1} \int_{-\infty}^{\infty} (x_2 - x_1) p(x_2; t_2 | x_1; t_1) dx_2 \\
&\quad + \frac{1}{\Delta t} \left. \frac{\partial^2 p}{\partial x_2^2} \right|_{x_2=x_1} \int_{-\infty}^{\infty} (x_2 - x_1)^2 p(x_2; t_2 | x_1; t_1) dx_2 + \frac{o\left[(x_2 - x_1)^2\right]}{\Delta t} = 0
\end{aligned}$$

Consider the $\lim t_2 \rightarrow t_1 \Rightarrow$

$$\frac{\partial p}{\partial t_1} + \alpha_1(x_1, t_1) \frac{\partial p}{\partial x_1} + \frac{1}{2} \alpha_2(x_1, t_1) \frac{\partial^2 p}{\partial x_1^2} = 0$$

$$p \equiv p(x_3, t_3 | x_1, t_1)$$

In the standard form the above equation written as

$$\frac{\partial p}{\partial t_0} = -\alpha_1(x_0, t_0) \frac{\partial p}{\partial x_0} - \frac{1}{2} \alpha_2(x_0, t_0) \frac{\partial^2 p}{\partial x_0^2}; p \equiv p(x, t | x_0, t_0)$$

is known as the backward Kolmogorov equation.

(Backward \because the independent variable $t_0 < t$)

$$\text{ICS: } p(x, t_0 | x_0, t_0) = \delta(x - x_0)$$

$$[\text{Typical}] \text{ BCS: } \lim_{x_0 \rightarrow \pm\infty} p(x, t | x_0, t_0) \rightarrow 0$$

Vector version of Backward Kolmogorov equation

$$dX(t) = f[t, X(t)]dt + G[t, X(t)]dB(t); t \geq 0; X(0) = X_0$$

$$\frac{\partial p}{\partial t_0} = -\sum_{j=1}^n f_j[t_0, \tilde{x}_0] \frac{\partial p}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n \left[GDG^t(t_0, \tilde{x}_0) \right]_{ij} \frac{\partial^2 p}{\partial x_{0i} \partial x_{0j}}$$

$$p \equiv p(\tilde{x}; t | \tilde{x}_0; t_0)$$

$$\text{ICS: } p(\tilde{x}; t_0 | \tilde{x}_0; t_0) = \prod_{i=1}^n \delta(x_i - x_{i0})$$

$$\text{BCS: } \forall j = 1, 2, \dots, n, \lim_{x_j \rightarrow \pm\infty} p(\tilde{x}; t | \tilde{x}_0; t_0) \rightarrow 0$$

Remarks

- If we are interested in pdf of $X(t)$ for a given initial condition we use forward equation.
- If we are interested in the pdf of time for first passage across a threshold, we use backward equation.

First passage times

Consider the scalar SDE

$$dX(t) = f(X)dt + G(X)dB(t) \text{ with } X(t_0) = x_0 < \underline{x_c}$$

T = time at which $X(t)$ reaches the critical value x_c for the first time given that $X(t_0) = x_0 < x_c$.

Safe region: $X(t) \leq x_c$

Unsafe region: $X(t) > x_c$

$R(t, x_c; x_0; t_0)$ = Probability that $X(t)$ stays in the safe region during the interval t_0 to t .

$$R(t, x_c; x_0; t_0) = \int_{x_l}^{x_c} p(x; t | x_0; t_0) dx = P[T > t - t_0 | X(t_0) = x_0]$$

$$R(t, x_c; x_0; t_0) = \int_{x_l}^{x_c} p(x; t | x_0; t_0) dx = P[T > t - t_0 | X(t_0) = x_0]$$

Remark

The samples of $X(t)$ are such that the first time they reach the level x_c they get absorbed. That is, $X(t) = x_c$ is an absorbing barrier.

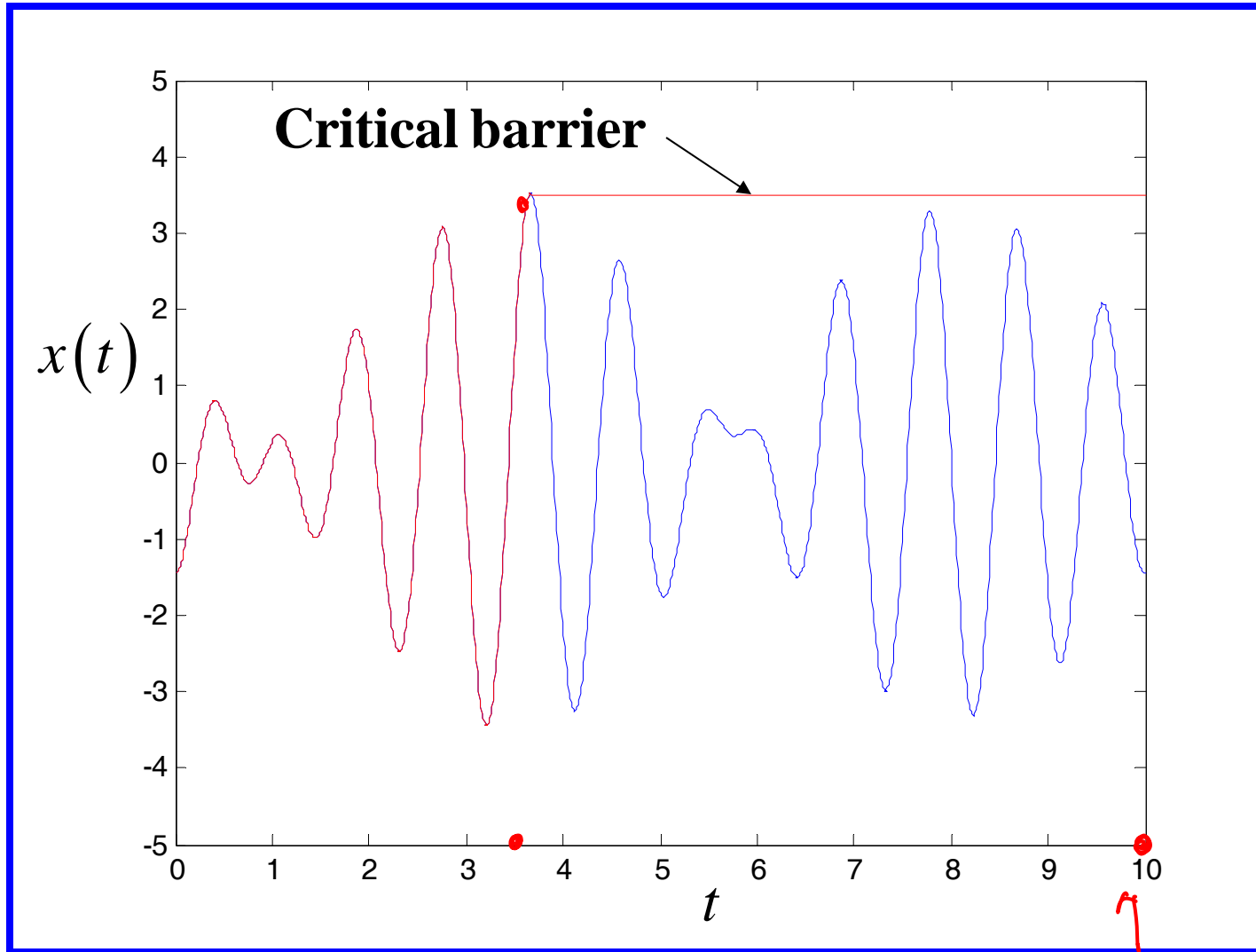
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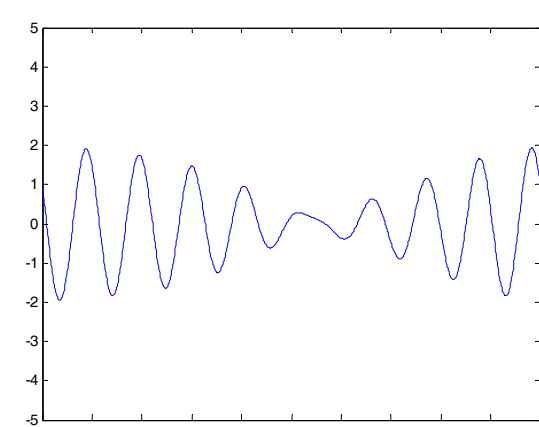
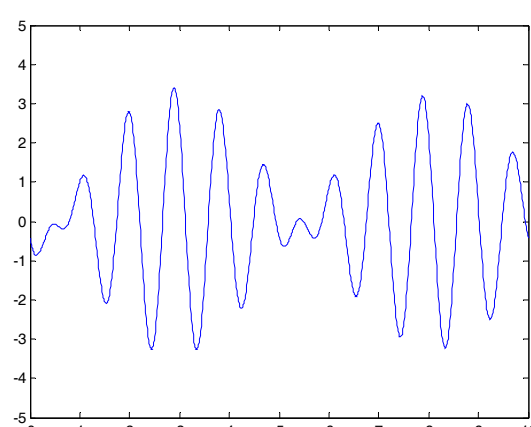
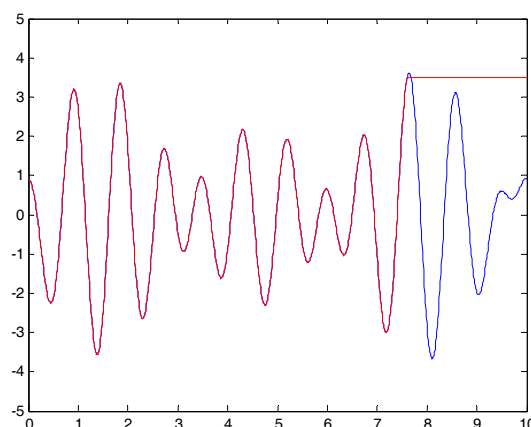
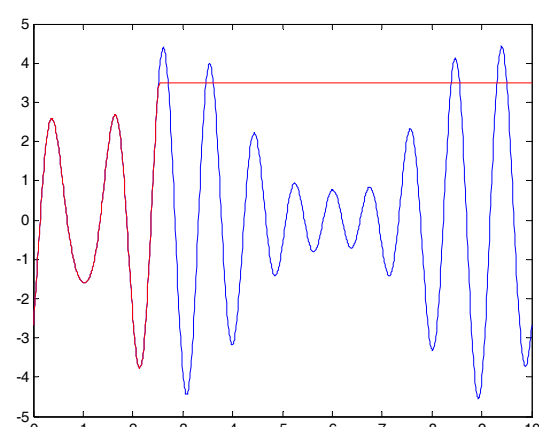
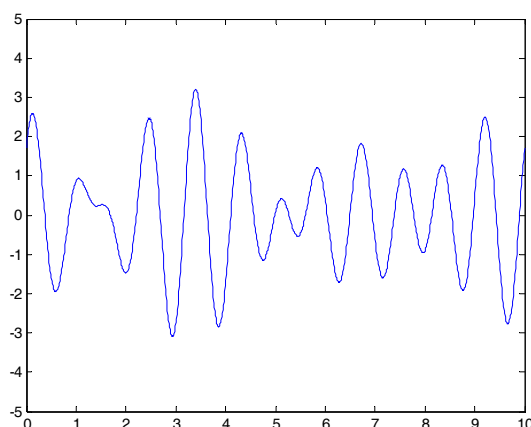
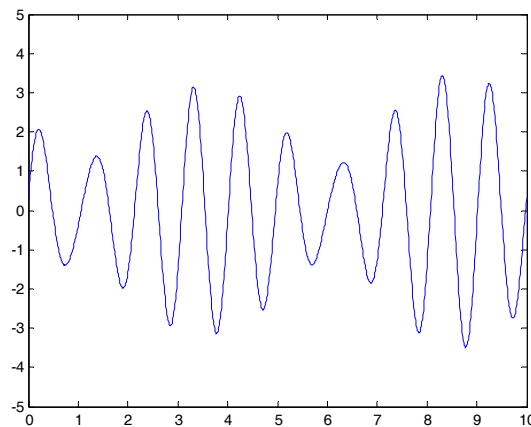
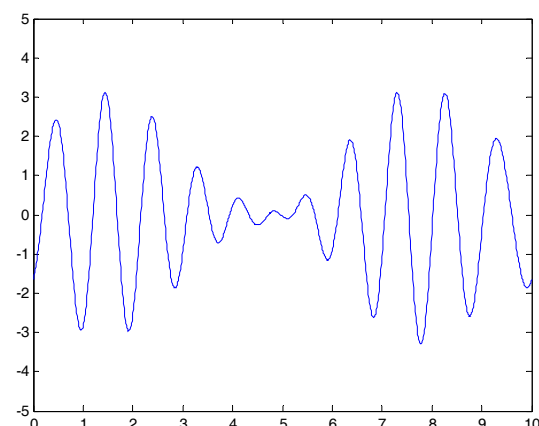
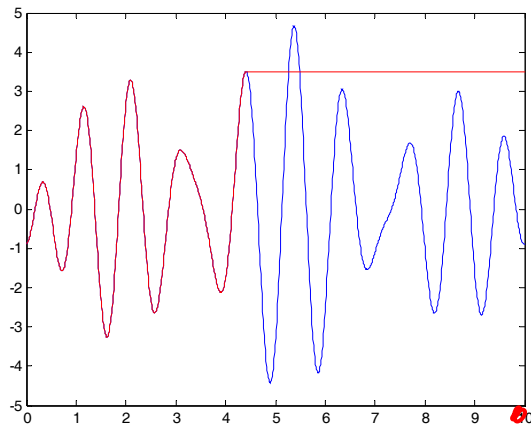
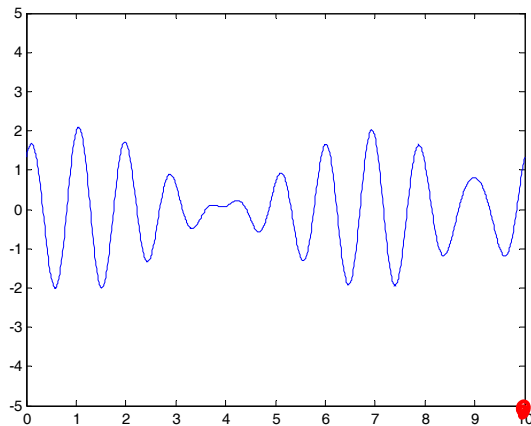
$$\frac{\partial p}{\partial t_0} + f(x_0) \frac{\partial p}{\partial x_0} + G(x_0) D \frac{\partial^2 p}{\partial x_0^2} = 0; p \equiv p(x; t | x_0; t_0)$$

$$\text{ICS: } p(x; t_0 | x_0; t_0) = \delta(x - x_0)$$

$$\text{BCS: } p(x_c; t | x_0; t_0) = 0 \quad [\text{Absorbing boundary}]$$

We consider only those trajectories which have not yet reached the critical barrier.





$$R(t, x_c; x_0; t_0) = \int_{x_l}^{x_c} p(x; t | x_0; t_0) dx = P[T > t - t_0 | X(t_0) = x_0]$$

$$\frac{\partial p}{\partial t_0} + f(x_0) \frac{\partial p}{\partial x_0} + G(x_0) D \frac{\partial^2 p}{\partial x_0^2} = 0 \Rightarrow$$

$$\frac{\partial R}{\partial t_0} + f(x_0) \frac{\partial R}{\partial x_0} + G(x_0) D \frac{\partial^2 R}{\partial x_0^2} = 0$$

BCS

$$R(t_0, x_c; x_0; t_0) = 1$$

- [At $t = t_0$, no trajectory has reached the critical barrier]

$$R(t, x_c; x_c; t_0) = 0$$

- [If the trajectory originates from x_c the probability of failure=1]

$$0 \leq R(t, x; x_0; t_0) \leq 1$$

- [Because R is a probability]

Summary

$$R(t, x_c; x_0; t_0) = P[T > t - t_0 \mid X(t_0) = x_0]$$

$$\frac{\partial R}{\partial t_0} + f(x_0) \frac{\partial R}{\partial x_0} + DG(x_0) \frac{\partial^2 R}{\partial x_0^2} = 0$$

$$R(t_0, x_c; x_0; t_0) = 1$$

$$R(t, x_c; x_c; t_0) = 0$$

$$0 \leq R(t, x; x_0; t_0) \leq 1$$

Assume: $x_0 = z_l$ is reflective

$$\left[\frac{\partial p(x; t \mid x_0; t_0)}{\partial x_0} \right]_{x_0 = z_l} = 0 \implies \left(\frac{\partial R}{\partial x_0} \right)_{x_0 = z_l} = 0$$

Let $\tau = t - t_0$

$$R(t, x_c; x_0; t_0) \equiv R(\tau, x_c; x_0) = P[T > \tau \mid X(t_0) = x_0]$$

$$-\frac{\partial R}{\partial \tau} + f(x_0) \frac{\partial R}{\partial x_0} + DG(x_0) \frac{\partial^2 R}{\partial x_0^2} = 0$$

$$R(\tau, x_c; x_0) |_{\tau=0} = 1$$

$$R(\tau, x_c; x_c) = 0$$

$$0 \leq R(\tau, x; x_0) \leq 1$$

$$\left(\frac{\partial R}{\partial x_0} \right)_{x_0=x_l} = 0$$

Moments of first passage time

$$P\left[T \leq \tau \mid X(t_0) = x_0\right] = 1 - P\left[T > \tau \mid X(t_0) = x_0\right]$$

$$P_T(\tau, x_c, x_0) = 1 - R(\tau, x_c; x_0)$$

$$p_T(\tau, x_c, x_0) = \frac{\partial}{\partial \tau} P_T(\tau, x_c, x_0) = -\frac{\partial}{\partial \tau} R(\tau, x_c; x_0)$$

We have

$$-\frac{\partial R}{\partial \tau} + f(x_0) \frac{\partial R}{\partial x_0} + DG(x_0) \frac{\partial^2 R}{\partial x_0^2} = 0$$

$$\Rightarrow p_T(\tau, x_c, x_0) = -f(x_0) \frac{\partial R}{\partial x_0} - DG(x_0) \frac{\partial^2 R}{\partial x_0^2}$$

$$\langle T^n \rangle = \int_0^{\infty} \tau^n p_T(\tau, x_c, x_0) d\tau = M_n(x_c, x_0)$$

$$\begin{aligned}
M_n &= \langle T^n \rangle = \int_0^\infty \tau^n p_T(\tau, x_c, x_0) d\tau = \int_0^\infty \tau^n \left(-\frac{\partial}{\partial \tau} R(\tau, x_c; x_0) \right) d\tau \\
&= \left[\tau^n R(\tau, x_c; x_0) \right]_0^\infty + n \int_0^\infty \tau^{n-1} R(\tau, x_c; x_0) d\tau \\
&= n \int_0^\infty \tau^{n-1} R(\tau, x_c; x_0) d\tau
\end{aligned}$$

We have

$$-\frac{\partial R}{\partial \tau} + f(x_0) \frac{\partial R}{\partial x_0} + DG(x_0) \frac{\partial^2 R}{\partial x_0^2} = 0$$

Multiply both sides of this equation by τ^n and integrate over 0 to $\infty \Rightarrow$

$$\int_0^{\infty} \tau^n \left(-\frac{\partial R}{\partial \tau} \right) d\tau + \int_0^{\infty} \tau^n f(x_0) \left(\frac{\partial R}{\partial x_0} \right) d\tau + D \int_0^{\infty} \tau^n G(x_0) \left(\frac{\partial^2 R}{\partial x_0^2} \right) d\tau = 0$$

$$M_n + \frac{1}{n+1} f(x_0) \frac{d}{dx_0} M_{n+1} + D \frac{G(x_0)}{n+1} \frac{d^2}{dx_0^2} M_{n+1} = 0$$

$$(n+1)M_n + f(x_0) \frac{d}{dx_0} M_{n+1} + DG(x_0) \frac{d^2}{dx_0^2} M_{n+1} = 0$$

$$M_n(x_c, x_0) \Big|_{x_0=x_c} = 0$$

$$M_n(x_c, x_0) \Big|_{x_0=x_l} < \infty$$

$$n = 1, 2, \dots$$

$$(n+1)M_n + f(x_0) \frac{d}{dx_0} M_{n+1} + DG(x_0) \frac{d^2}{dx_0^2} M_{n+1} = 0$$

$$n = 0 \Rightarrow M_0 = 1$$

$$1 + f(x_0) \frac{dM_1}{dx_0} + DG(x_0) \frac{d^2 M_1}{dx_0^2} = 0$$

$$M_1(x_c, x_0) \big|_{x_0=x_c} = 0$$

$$M_1(x_c, x_0) \big|_{x_0=x_l} < \infty$$

$$n = 1$$

$$f(x_0) \frac{dM_2}{dx_0} + DG(x_0) \frac{d^2 M_2}{dx_0^2} = -2M_1$$

$$M_2(x_c, x_0) \big|_{x_0=x_c} = 0$$

$$M_2(x_c, x_0) \big|_{x_0=x_l} < \infty$$

VitH GPR eqns

Generalized Pontryagin Equations

$$f(x_0) \frac{d}{dx_0} M_{n+1} + DG(x_0) \frac{d^2}{dx_0^2} M_{n+1} = -(n+1) M_n$$

$$M_{n+1}(x_c, x_0) \Big|_{x_0=x_c} = 0$$

$$M_{n+1}(x_c, x_0) \Big|_{x_0=x_l} < \infty$$

$$n = 1, 2, \dots$$

Generalization to multi - dimensional problems

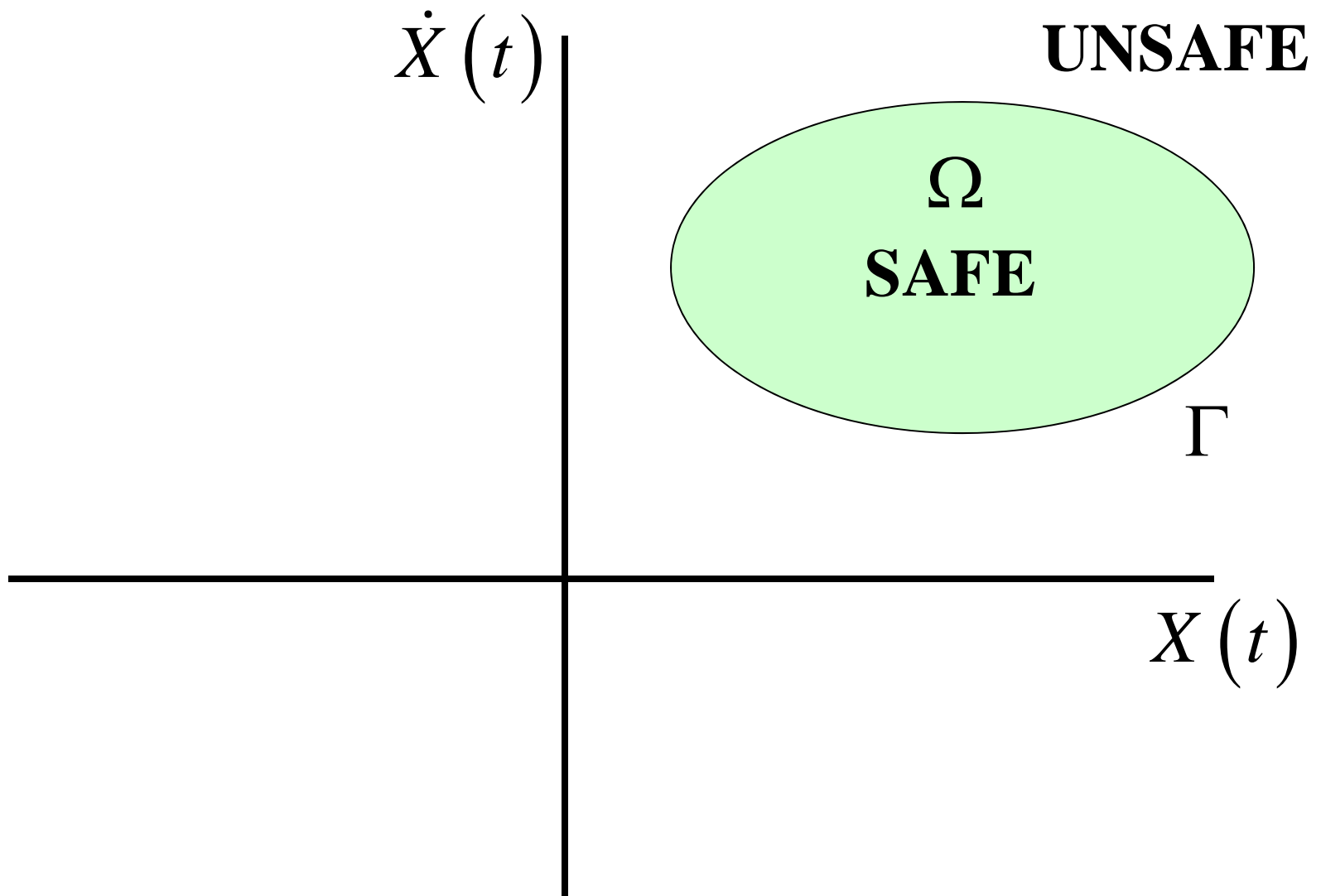
$$dX(t) = f[t, X(t)]dt + G[t, X(t)]dB(t); t \geq 0; X(0) = X_0$$

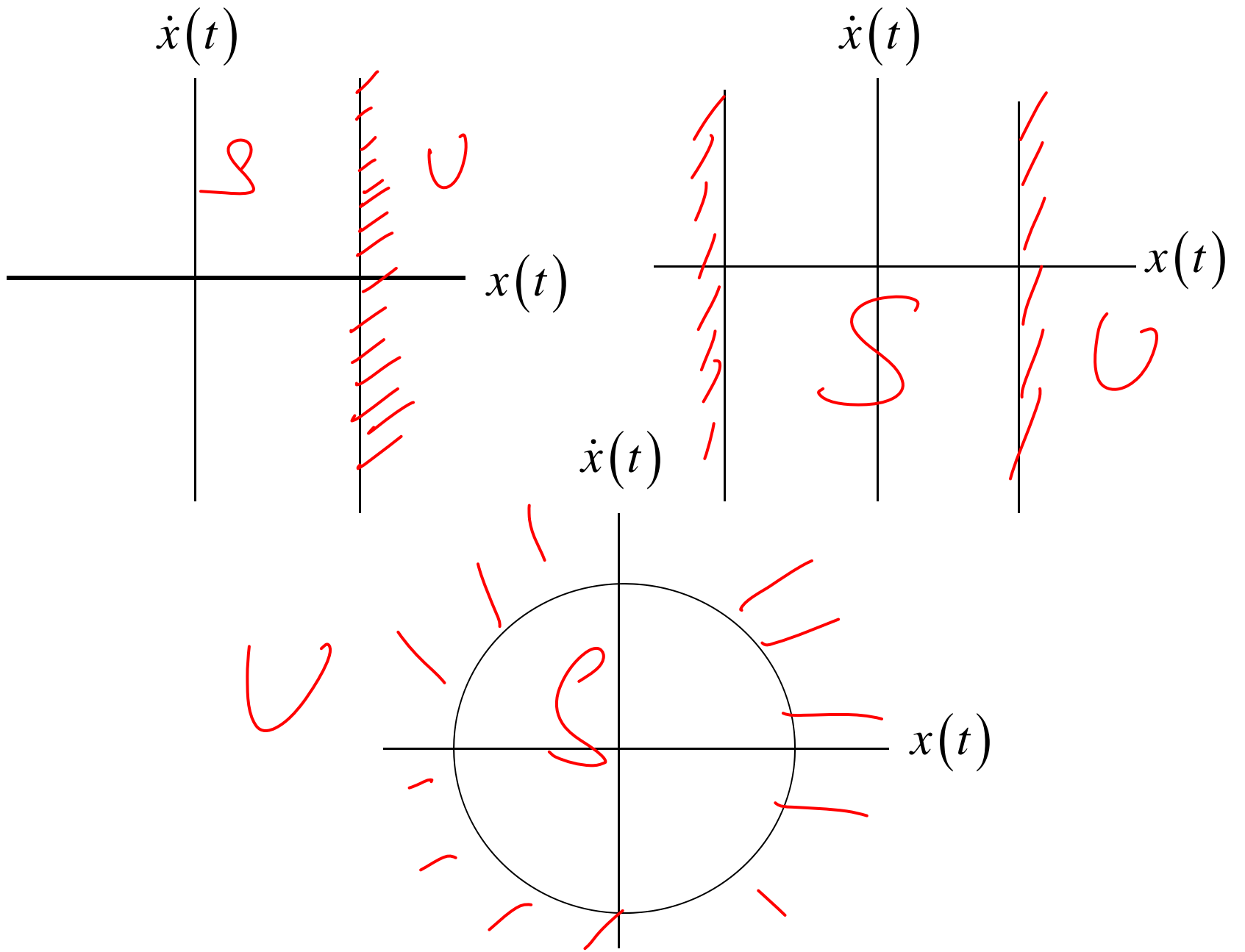
$$\frac{\partial p}{\partial t_0} = -\sum_{j=1}^n f_j[t_0, \tilde{x}_0] \frac{\partial p}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n [GDG^t(t_0, \tilde{x}_0)]_{ij} \frac{\partial^2 p}{\partial x_{0i} \partial x_{0j}}$$

$$p \equiv p(\tilde{x}; t | \tilde{x}_0; t_0)$$

Safe region : $X(t) \in \Omega$; Γ = Limits of safe region

$$R(t, \Gamma; t_0, \Omega) = \int_{\Omega} p(\tilde{x}; t | \tilde{x}_0; t_0) d\tilde{x} = P[T > t - t_0 | X(t_0) = \tilde{x}_0]$$





$$R(t, \Gamma; t_0, \Omega) = \int_{\Omega} p(\tilde{x}; t | \tilde{x}_0; t_0) d\tilde{x} = P[T > t - t_0 | X(t_0) = \tilde{x}_0]$$

$$\frac{\partial R}{\partial t_0} = - \sum_{j=1}^n f_j[t_0, \tilde{x}_0] \frac{\partial R}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n [GDG^t(t_0, \tilde{x}_0)]_{ij} \frac{\partial^2 R}{\partial x_{0i} \partial x_{0j}}$$

$$R(t_0, \Gamma; t_0, \Omega) = 1$$

- [At $t = t_0$, no trajectory has reached the critical barrier]

$$R(t, \Gamma; t_0, \Gamma) = 0$$

- [If the trajectory originates on Γ the probability of failure=1]

$$0 \leq R(t, \Gamma; t_0, \Omega) \leq 1$$

- [Because R is a probability]

Let $\tau = t - t_0$

$$R(t, \Gamma; x_0; t_0) \equiv R(\tau, \Gamma; x_0) = P[T > \tau \mid X(t_0) = x_0 \in \Omega]$$

$$-\frac{\partial R}{\partial \tau} = -\sum_{j=1}^n f_j[t_0, \tilde{x}_0] \frac{\partial R}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n [GDG^t(t_0, \tilde{x}_0)]_{ij} \frac{\partial^2 R}{\partial x_{0i} \partial x_{0j}}$$

$$R(\tau, \Gamma; x_0 \in \Omega) \big|_{\tau=0} = 1$$

$$R(\tau, \Gamma; x_c \in \Gamma) = 0$$

$$0 \leq R(\tau, \Gamma; x_0 \in \Omega) \leq 1$$

Moments of first passage time

$$P\left[T \leq \tau \mid X(t_0) = x_0 \in \Omega\right] = 1 - P\left[T > \tau \mid X(t_0) = x_0 \in \Omega\right]$$

$$R(t, \Gamma; \tilde{x}_0 \in \Omega) = \int_{\Omega} p(\tilde{x}; t \mid \tilde{x}_0; t_0) d\tilde{x} = P\left[T > t - t_0 \mid X(t_0) = \tilde{x}_0 \in \Omega\right]$$

$$P_T(\tau, \Gamma, x_0 \in \Omega) = 1 - R(t, \Gamma; \tilde{x}_0 \in \Omega)$$

$$p_T(\tau, \Gamma, x_0 \in \Omega) = \frac{\partial}{\partial \tau} P_T(\tau, \Gamma, x_0 \in \Omega) = -\frac{\partial}{\partial \tau} R(t, \Gamma; \tilde{x}_0 \in \Omega)$$

We have

$$-\frac{\partial R}{\partial \tau} = -\sum_{j=1}^n f_j [t_0, \tilde{x}_0] \frac{\partial R}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n \left[G D G^t (t_0, \tilde{x}_0) \right]_{ij} \frac{\partial^2 R}{\partial x_{0i} \partial x_{0j}}$$

$$\Rightarrow p_T (\tau, \Gamma, x_0 \in \Omega) = -\sum_{j=1}^n f_j [t_0, \tilde{x}_0] \frac{\partial R}{\partial x_{0j}} -$$

$$\sum_{i=1}^n \sum_{j=1}^n \left[G D G^t (t_0, \tilde{x}_0) \right]_{ij} \frac{\partial^2 R}{\partial x_{0i} \partial x_{0j}}$$

$$\langle T^n \rangle = \int_0^{\infty} \tau^n p_T (\tau, \Gamma, x_0 \in \Omega) d\tau = M_n (\Gamma, x_0 \in \Omega)$$

\Rightarrow

$$\mathbf{L} = -\sum_{j=1}^n f_j [t_0, \tilde{x}_0] \frac{\partial}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n \left[G D G^t (t_0, \tilde{x}_0) \right]_{ij} \frac{\partial^2}{\partial x_{0i} \partial x_{0j}}$$

$$(n+1) M_n + \mathbf{L} M_{n+1} = 0; n = 1, 2, \dots$$

GTV

Generalized Pontryagin Equations

$$(n+1)M_n + \mathbf{L}M_{n+1} = 0; n = 1, 2, \dots$$

$$\mathbf{L} = -\sum_{j=1}^n f_j [t_0, \tilde{x}_0] \frac{\partial}{\partial x_{0j}} - \sum_{i=1}^n \sum_{j=1}^n \left[G D G^t (t_0, \tilde{x}_0) \right]_{ij} \frac{\partial^2}{\partial x_{0i} \partial x_{0j}}$$

$$M_{n+1} (\Gamma, x_0 \in \Gamma) = 0$$

$$M_{n+1} (\Gamma, x_0 \in \Omega) < \infty$$

$$n = 1, 2, \dots$$

Introduction to method of stochastic averaging

- An approximate method for analysing lightly damped nonlinear systems under broad band excitations.
- Leads to simplified models for response envelope and phase processes
- Enables application of Markovian methods to the solution of equations governing the response envelope and phase processes
- Best suited to study nonlinear sdof systems.

Reference :

J B Roberts and P D Spanos, 1986, Stochastic averaging: an approximate method of solving nonlinear random vibration problems, International Journal of Nonlinear Mechanics, 21(2),111-134.

Illustration of deterministic averaging procedure

Consider the free vibration response of a nonlinearly damped sdof system

$$\ddot{u} + \omega_0^2 u = \varepsilon f(u, \dot{u}); t \geq 0; u(0) \& \dot{u}(0) \text{ specified.}$$

$\varepsilon =$ a small non-dimensional number

$$\varepsilon = 0 \Rightarrow$$

$$\left. \begin{aligned} u(t) &= a \cos(\omega_0 t + \beta) \\ \dot{u}(t) &= -a\omega_0 \sin(\omega_0 t + \beta) \end{aligned} \right\}$$

For $\varepsilon \neq 0$, we consider the transformation

$$\begin{aligned} u(t) &= a(t) \cos[\omega_0 t + \beta(t)] \\ \dot{u}(t) &= -a(t) \omega_0 \sin[\omega_0 t + \beta(t)] \end{aligned}$$

Note: this step involves no approximations

$$\ddot{u} + \omega_0^2 u = 0$$

$$u(t) = A \cos \omega_0 t + B \sin \omega_0 t$$

$$u(0) = u_0 \Rightarrow A = u_0$$

$$\dot{u}(0) = \dot{u}_0 \Rightarrow B = \frac{\dot{u}_0}{\omega_0}$$

$$\begin{aligned} u(t) &= u_0 \cos \omega_0 t + \frac{\dot{u}_0}{\omega_0} \sin \omega_0 t \\ &= R \cos(\omega_0 t + \beta) \end{aligned}$$

Consider

$$u(t) = a(t) \cos [\omega_0 t + \beta(t)] = a(t) \cos \phi(t); \phi(t) = \omega_0 t + \beta(t)$$

$$\Rightarrow \dot{u}(t) = \dot{a}(t) \cos \phi(t) - a(t) [\omega_0 + \dot{\beta}(t)] \sin \phi(t)$$

Since we are taking $\dot{u}(t) = -a(t) \omega_0 \sin \phi(t)$, it follows

$$\dot{a}(t) \cos \phi(t) - a(t) \dot{\beta}(t) \sin \phi(t) = 0 \dots (A)$$

Consider $\dot{u}(t) = -a(t) \omega_0 \sin \phi(t)$

$$\Rightarrow \ddot{u}(t) = -\dot{a}(t) \omega_0 \sin \phi(t) - a(t) \omega_0 [\omega_0 + \dot{\beta}(t)] \cos \phi(t)$$

$$\Rightarrow -\dot{a}(t) \omega_0 \sin \phi(t) - a(t) \omega_0 [\omega_0 + \dot{\beta}(t)] \cos \phi(t) +$$

$$\omega_0^2 a(t) \cos \phi(t) = \varepsilon f [a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t)]$$

$$\Rightarrow -\dot{a}(t) \omega_0 \sin \phi(t) - a(t) \omega_0 \cos \phi(t) =$$

$$\varepsilon f [a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t)] \dots (B)$$

$$\dot{a}(t) \cos \phi(t) - a(t) \dot{\beta}(t) \sin \phi(t) = 0 \dots (A)$$

$$-\dot{a}(t) \omega_0 \sin \phi(t) - a(t) \omega_0 \cos \phi(t) =$$

$$\varepsilon f \left[a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t) \right] \dots (B)$$

Solving for $\dot{a}(t)$ and $\dot{\beta}(t)$ we get

$$\dot{a}(t) = -\frac{\varepsilon}{\omega_0} \sin \phi(t) f \left[a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t) \right]$$

$$\dot{\beta}(t) = \frac{\varepsilon}{\omega_0 a(t)} \cos \phi(t) f \left[a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t) \right]$$

- Note: no approximations till now.
- If ε is small, $\dot{a}(t)$ and $\dot{\beta}(t)$ are small.

That is $a(t)$ and $\beta(t)$ do not change appreciably over

the time duration t to $t + \frac{2\pi}{\omega_0}$.

$$\dot{a}(t) = -\frac{\varepsilon}{\omega_0} \sin \phi(t) f \left[a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t) \right]$$

$$\dot{\beta}(t) = \frac{\varepsilon}{\omega_0 a(t)} \cos \phi(t) f \left[a(t) \cos \phi(t), -a(t) \omega_0 \sin \phi(t) \right]$$

$$\dot{a}(t) \approx -\frac{\varepsilon}{\omega_0} \frac{1}{2\pi} \int_0^{2\pi} \sin \phi f \left[a \cos \phi, -a \omega_0 \sin \phi \right] d\phi$$

$$\dot{\beta}(t) \approx -\frac{\varepsilon}{\omega_0 a} \frac{1}{2\pi} \int_0^{2\pi} \cos \phi f \left[a \cos \phi, -a \omega_0 \sin \phi \right] d\phi$$

These are averaged equations.

Example

$$\ddot{u} + \omega_0^2 u = -2\varepsilon\mu\dot{u}$$

$$\dot{a}(t) \approx -\frac{\varepsilon}{\omega_0} \frac{1}{2\pi} \int_0^{2\pi} \sin \phi [2\mu a \omega_0 \sin \phi] d\phi$$

$$= -\frac{\varepsilon}{\omega_0} \frac{1}{2\pi} (2\mu a \omega_0) \int_0^{2\pi} \sin^2 \phi d\phi = -\varepsilon a \mu$$

$$\dot{\beta}(t) \approx -\frac{\varepsilon}{\omega_0 a} \frac{1}{2\pi} \int_0^{2\pi} \cos \phi [2\mu a \omega_0 \sin \phi] d\phi = 0$$

$$\Rightarrow u(t) \approx a_0 \exp(-\varepsilon\mu t) \cos[\omega_0 t + \beta_0] //$$

Exercise: compare this with the known exact solution

Example

$$\ddot{u} + \omega_0^2 u = -\varepsilon \dot{u} |\dot{u}|$$

$$\dot{a}(t) \approx -\frac{\varepsilon}{\omega_0} \frac{1}{2\pi} \int_0^{2\pi} \sin \phi \left[a\omega_0 \sin \phi | -a\omega_0 \sin \phi | \right] d\phi$$

$$= -\frac{\varepsilon}{\omega_0} \frac{1}{2\pi} \left(a^2 \omega_0^2 \right) \int_0^{2\pi} \sin^2 \phi |\sin \phi| d\phi = -\varepsilon \frac{4\omega_0 a^2}{3\pi}$$

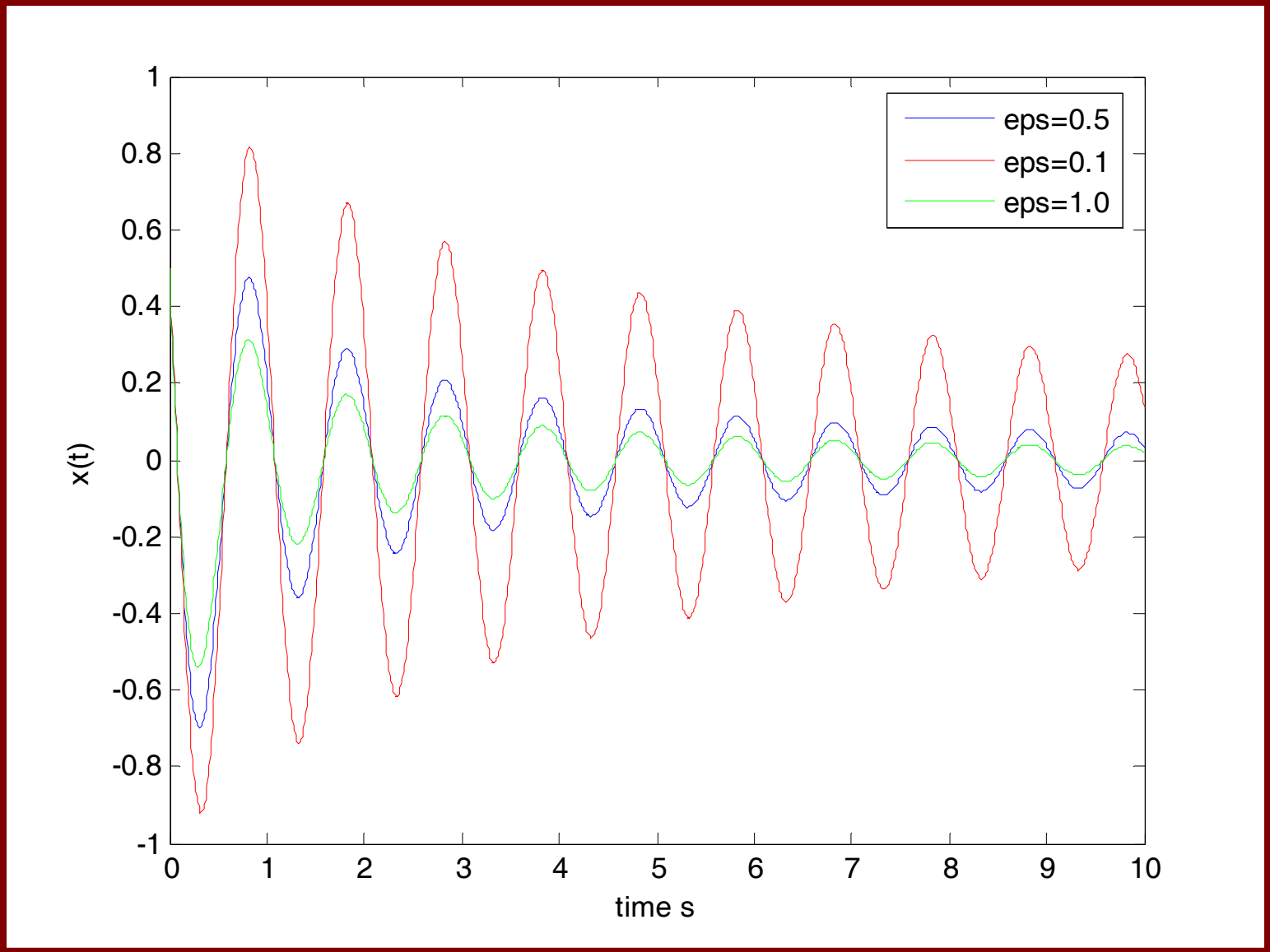
$$\dot{\beta}(t) \approx -\frac{\varepsilon}{\omega_0 a} \frac{1}{2\pi} \int_0^{2\pi} \cos \phi \left[a\omega_0 \sin \phi | -a\omega_0 \sin \phi | \right] d\phi = 0$$

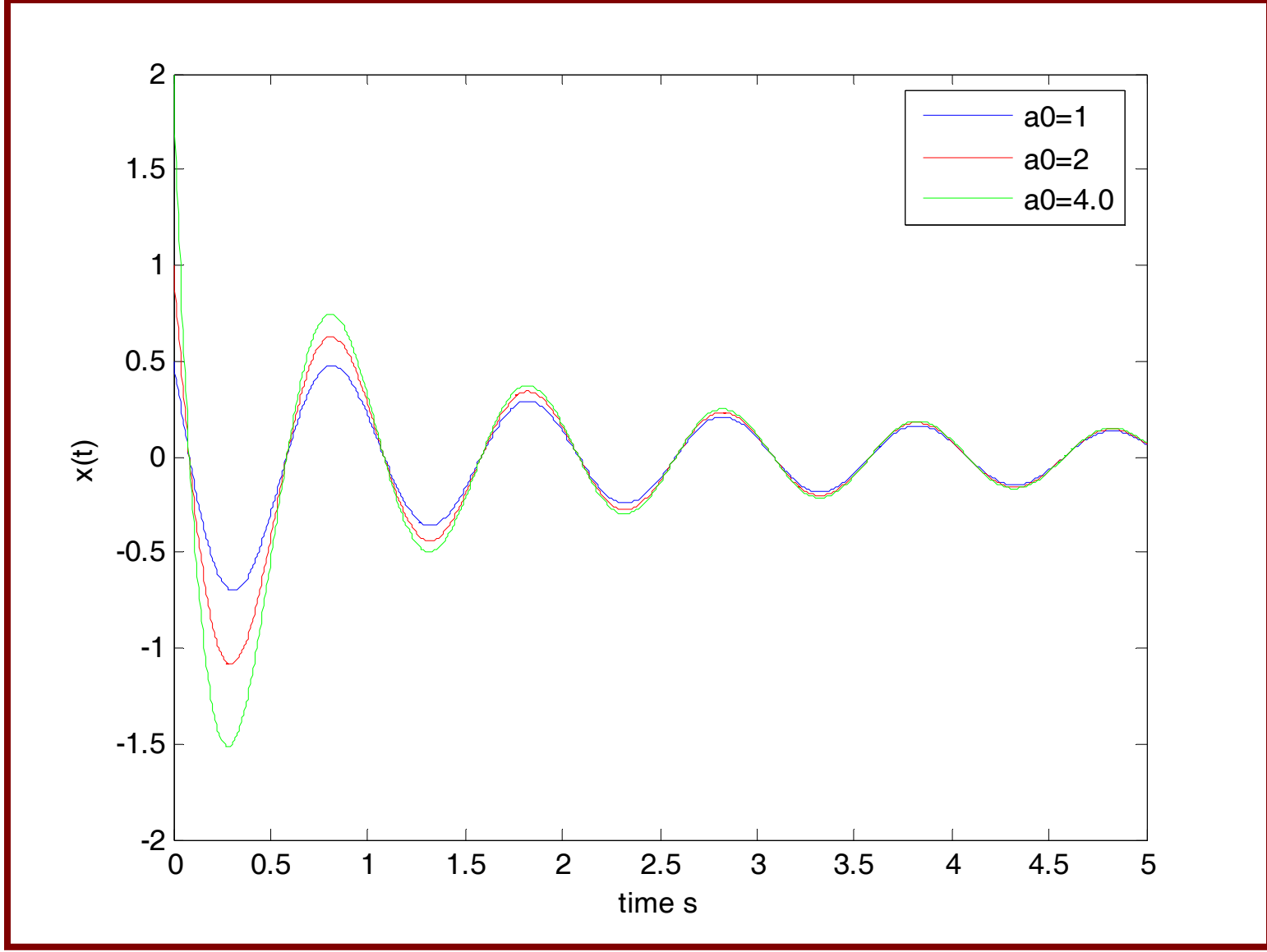
$$\Rightarrow u(t) \approx \left\{ \frac{a_0}{1 + \frac{4\varepsilon\omega_0 a_0}{3\pi} t} \right\} \cos[\omega_0 t + \beta_0]$$

$$\dot{a} = -\varepsilon \frac{4\omega_0^2 a^2}{3\pi}$$

$$\frac{da}{a^2} = -\varepsilon \frac{dt 4\omega_0^2}{3\pi}$$

Note: Rate of free vibration decay depends on ics.





Example

$$\ddot{u} + \omega_0^2 u = \varepsilon (\dot{u} - \dot{u}^3) -$$

$$\ddot{u} + \omega_0^2 u - \varepsilon \dot{u} (1 - \dot{u}^2) = 0$$

$$\dot{a} = \frac{1}{2} \varepsilon a \left[1 - \frac{3}{4} \omega_0^2 a^2 \right] \& \dot{\beta} = 0$$

$$a(t) = \left\{ \frac{a_0^2}{\frac{3}{4} \omega_0^2 a_0^2 + \left[1 - \frac{3}{4} \omega_0^2 a_0^2 \right] \exp(-\varepsilon t)} \right\}^{\frac{1}{2}} \& \underline{\underline{\beta = \beta_0}}$$

$$\Rightarrow u(t) \approx \left\{ \frac{a_0^2}{\frac{3}{4} \omega_0^2 a_0^2 + \left[1 - \frac{3}{4} \omega_0^2 a_0^2 \right] \exp(-\varepsilon t)} \right\}^{\frac{1}{2}} \cos[\omega_0 t + \beta_0]$$

$$u(t) \approx \left\{ \frac{a_0^2}{\frac{3}{4} \omega_0^2 a_0^2 + \left[1 - \frac{3}{4} \omega_0^2 a^2 \right] \exp(-\varepsilon t)} \right\}^{\frac{1}{2}} \cos[\omega_0 t + \beta_0]$$

$$\lim_{t \rightarrow \infty} u(t) \rightarrow \left\{ \frac{4}{3\omega_0^2} \right\}^{\frac{1}{2}} \cos[\omega_0 t + \beta_0]$$

Note: the amplitude of steady state oscillations (in free vibrations) is independent of initial conditions. Contrast this with the response of undamped linear sdof system.

