Stochastic Structural Dynamics

Lecture-34

Probabilistic methods in earthquake engineering-3

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Examples of stochastic models for earthquake ground motions

- Single component: stationary & nonstationary models
- Multi-component and spatially varying load models
- •Gaussian and Poisson pulse process models

Main concerns

- frequency content
- \bullet transient nature and duration
- time dependent frequency content
- multi-component nature
- spatial variability
- \bullet translations and rotations
- models for displacement and velocity components
- seismological considerations

Spectral representation of an evolutionary random process
\n
$$
X(t) = \int_{-\infty}^{\infty} a(t, \omega) \exp(i\omega t) dZ(\omega)
$$
\n
$$
a(t, \omega) = \text{deterministic function (in general, complex valued)}
$$
\n
$$
Z(\omega) = \text{orthogonal increment random process (complex valued)}
$$
\nwith $\langle dZ(\omega) \rangle = 0 \& \langle dZ(\omega_1) dZ^*(\omega_2) \rangle = \delta(\omega_1 - \omega_2) d\Psi(\omega)$ \n
$$
\langle X(t_1) X^*(t_2) \rangle = \int_{-\infty}^{\infty} a(t_1, \omega) a^*(t_2, \omega) \exp[i\omega(t_1 - t_2)] d\Psi(\omega)
$$
\n
$$
\sigma_X^2(t) = \int_{-\infty}^{\infty} |a(t, \omega)|^2 d\Psi(\omega)
$$
\nIf $d\Psi(\omega) = \Phi(\omega) d\omega$, we get $\sigma_X^2(t) = \int_{-\infty}^{\infty} |a(t, \omega)|^2 \Phi(\omega) d\omega$
\nWe interpret $S_{XX}(\omega) = |a(t, \omega)|^2 \Phi(\omega)$ as the nonstationary (evolutionary) PSD function of $X(t)$.

3

Filtered Poisson Process models for earthquake ground motions Rationale

During earthquakes slips occur along fault lines in an intermittent manner. This sends out a train of stress waves in the earth cr ust. This eventually results in ground shaking. **Recall**

$$
X(t) = \sum_{j=1}^{N(T)} Y_j w(t, \tau_j); 0 < t \le T
$$

$$
N(T) = \text{counting process, Poisson; arrival rate } = \lambda(t)
$$

$$
\tau_j = \text{arrival times; random}
$$

$$
w(t, \tau_j) = \text{Deterministic pulse shape } (=0 \forall t \le \tau_j).
$$

$$
Y_j = \text{random magnitude of the } j\text{-th pulse.}
$$

$$
m_X(t) = m_Y \int_0^t w(t, \tau) \lambda(\tau) d\tau,
$$

\n
$$
C_{XX}(t_1, t_2) = E(Y^2) \int_0^{\min(t_1, t_2)} w(t_1, \tau) w(t_2, \tau) \lambda(\tau) d\tau,
$$

\n
$$
\sigma_X^2(t) = E(Y^2) \int_0^t w^2(t, \tau) \lambda(\tau) d\tau
$$

Reference

Y K Lin and G C Cai, 1995, McGraw Hill, NY.

Let
$$
X(t) = \sum_{j=1}^{N(T)} Y_j w(t - \tau_j)
$$
; $0 < t \le T$
\n
$$
m_X(t) = m_Y \int_0^t w(t - \tau) \lambda(\tau) d\tau,
$$
\n
$$
C_{XX}(t_1, t_2) = E(Y^2) \int_0^{\min(t_1, t_2)} w(t_1 - \tau) w(t_2 - \tau) \lambda(\tau) d\tau,
$$
\nLet $\lambda(\tau) = 0 \forall \tau < 0$. Since $w(t - \tau) = 0 \forall t - \tau < 0$ we can write\n
$$
C_{XX}(t_1, t_2) = E(Y^2) \int_{-\infty}^{\infty} w(t_1 - \tau) w(t_2 - \tau) \lambda(\tau) d\tau
$$

We introduce
\n
$$
b(t, \omega) = \int_{-\infty}^{\infty} w(u) \sqrt{\lambda(t-u)} \exp(-i\omega u) du
$$
\nso that
\n
$$
w(u) \sqrt{\lambda(t-u)} = \frac{1}{2\pi} \int_{-\infty}^{\infty} b(t, \omega) \exp(i\omega u) d\omega
$$
\nLet $t-u = \tau \Rightarrow$
\n
$$
w(t-\tau) \sqrt{\lambda(\tau)} = \frac{1}{2\pi} \int_{-\infty}^{\infty} b(t, \omega) \exp[i\omega(t-\tau)] d\omega
$$
\nLHS is real \Rightarrow
\n
$$
w(t-\tau) \sqrt{\lambda(\tau)} = \frac{1}{2\pi} \int_{-\infty}^{\infty} b^*(t, \omega) \exp[-i\omega(t-\tau)] d\omega
$$

Substitute
\n
$$
w(t_1 - \tau)\sqrt{\lambda(\tau)} = \frac{1}{2\pi} \int_{-\infty}^{\infty} b(t_1, \omega) \exp[i\omega(t_1 - \tau)] d\omega
$$
\n
$$
w(t_2 - \tau)\sqrt{\lambda(\tau)} = \frac{1}{2\pi} \int_{-\infty}^{\infty} b^*(t_2, \omega) \exp[-i\omega(t_2 - \tau)] d\omega
$$
\ninto $C_{XX}(t_1, t_2) = E(Y^2) \int_{-\infty}^{\infty} w(t_1 - \tau) w(t_2 - \tau) \lambda(\tau) d\tau$
\nand noting that $\frac{1}{2\pi} \int_{-\infty}^{\infty} \exp[-i(\omega_2 - \omega_1)\tau] d\tau = \delta(\omega_2 - \omega_1) \Rightarrow$
\n $C_{XX}(t_1, t_2) = \frac{E(Y^2)}{2\pi} \int_{-\infty}^{\infty} b(t_1, \omega) b^*(t_2, \omega) \exp[-i\omega(t_2 - t_1)] d\omega$

$$
C_{XX}(t_1, t_2) = \frac{E(Y^2)}{2\pi} \int_{-\infty}^{\infty} b(t_1, \omega) b^*(t_2, \omega) \exp[-i\omega(t_2 - t_1)] d\omega
$$

\n
$$
\Rightarrow \sigma_X^2(t) = \frac{E(Y^2)}{2\pi} \int_{-\infty}^{\infty} |b(t, \omega)|^2 d\omega
$$

\n
$$
\Rightarrow S_{XX}(t, \omega) = \frac{E(Y^2)}{2\pi} |b(t, \omega)|^2
$$

\nwith
\n
$$
b(t, \omega) = \int_{-\infty}^{\infty} w(u) \sqrt{\lambda(t - u)} \exp(-i\omega u) du
$$

Selection of the shape of the pulse

Model -1

As in Kanai Tajimi model, the soil layer is modeled as an elastic half-space which can be represented as a sdof system.

$$
\begin{aligned}\n\ddot{u} + 2\eta_g \omega_g \dot{u} + \omega_g^2 u &= 2\eta_g \omega_g \dot{R} + \omega_g^2 R \\
H_1(\omega) &= \frac{\omega_g^2 + i2\eta_g \omega_g \omega}{\left(\omega_g^2 - \omega^2\right)^2 + \left(2\eta_g \omega_g \omega\right)^2} \\
h_1(t) &= \omega_g \exp\left(-\eta_g \omega_g t\right) \left\{\frac{1 - 2\eta_g^2}{\sqrt{1 - \eta_g^2}} \sin \omega_{gd} t + 2\eta_g \cos \omega_{gd} t\right\}; t > 0 \\
G(t) &= \sum_{j=1}^{N(T)} Y_j h_1(t - \tau_j)\n\end{aligned}
$$

Seismic wave amplification through soil layers

$$
u(z,t) = \phi(z) \exp(i\omega t)
$$

\n
$$
\Rightarrow -\rho \omega^2 \phi \exp(i\omega t) = G\phi'' \exp(i\omega t) + i\eta \omega \phi'' \exp(i\omega t)
$$

\n
$$
\Rightarrow \phi''(G + i\eta \omega) + \rho \omega^2 \phi = 0
$$

\n
$$
\Rightarrow \phi'' + \lambda^2 \phi = 0; \quad \lambda^2 = \frac{\rho \omega^2}{(G + i\eta \omega)}
$$

\n
$$
\phi(z) = A \cos \lambda z + B \sin \lambda z
$$

\n
$$
\phi(0) = 1 \quad \phi'(L) = 0
$$

\n
$$
\Rightarrow \phi(x) = \cos \lambda z + \tan \lambda L \sin \lambda z
$$

$$
\phi(L) = \frac{1}{\cos \lambda L} = \frac{1}{\cos \left(\frac{\omega L}{\nu}\right)} / \sqrt{\frac{1}{\nu^*}} = \sqrt{\frac{G(1 + i\omega \eta)}{\rho}} = \sqrt{\frac{G(1 + 2i\xi)}{\rho}}
$$

Selectron of the shape of the pulse
\nModel -2
\nSoil layer modeled as a shear beam with
\nhysteretic damping
\n
$$
\frac{\partial^2 w}{\partial t^2} - \beta^2 \frac{\partial^2 w}{\partial y^2} = 0
$$
\n
$$
H_2(\omega) = \left[\cos \left\{ \frac{\omega l}{\beta (1 + i \gamma \operatorname{sgn} \omega)} \right\} \right]^{-1}
$$
\n
$$
h_2(t) = \frac{2\beta}{l} \sum_{n=0}^{\infty} (-1)^n \exp \left[-\left(n + \frac{1}{2} \right) \frac{\pi \beta \gamma}{l} t \right]
$$
\n
$$
\left\{ \gamma \cos \left[\left(n + \frac{1}{2} \right) \frac{\pi \beta \gamma}{l} t \right] + \sin \left[\left(n + \frac{1}{2} \right) \frac{\pi \beta \gamma}{l} t \right] \right\} t > 0
$$
\n
$$
G(t) = \sum_{j=1}^{N(T)} Y_j h_2(t - \tau_j)
$$

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Selection of the shape of the pulse

Model -3

Soil layer modeled as a viscously damped shear beam

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$$
\frac{\partial^2 w}{\partial t^2} + \frac{1}{\tau_r} \frac{\partial w}{\partial t} - \beta^2 \frac{\partial^2 w}{\partial y^2} = 0
$$

\n
$$
H_3(\omega) = \left[\cos \left(\frac{l}{\beta} \sqrt{\omega^2 - i \frac{\omega}{\tau_r}} \right) \right]^{-1}
$$

\n
$$
h_3(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H_3(\omega) \exp(i \omega t) d\omega; \ t > 0
$$

\n
$$
G(t) = \sum_{j=1}^{N(T)} Y_j h_3(t - \tau_j)
$$

Selection of the shape of the pulse

Model -4

Soil layer modeled as a inhomogeneous hysteretically damped shear beam

$$
\frac{\partial^2 w}{\partial t^2} - \beta^2 \frac{\partial^2 w}{\partial y^2} - \beta^2 \frac{d \{\ln A(y)\}\partial w}{dy} = 0
$$

\n
$$
H_4(\omega) = \exp(-my) \left[\cos(\delta l) - \frac{m}{\delta} \sin(\delta l) \right]
$$

\n
$$
\delta = \left[\omega^2 \beta^{-2} (1 + iy \operatorname{sgn} \omega)^{-2} - m^2 \right]_4^{0.5}
$$

\n
$$
h_4(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H_3(\omega) \exp(i\omega t) d\omega; \ t > 0
$$

\n
$$
G(t) = \sum_{j=1}^{N(T)} Y_j h_4(t - \tau_j)
$$

 $\left(\mathbf{r},t\right)=\sum Y_{j}\mathcal{g}_{k}\left(\pmb{\kappa},t;\rho,\upsilon\right)$ $\big(\mathsf{\Gamma},t;\rho,v\big)$ 1 General model Green's function which describes the ground acceleration in the k – th direction at a site location^o and time t due to an impulsive application of a doubl e couple. *N(T)* $k \left(\cdot \right)$, $\left(\cdot \right)$ $\left(\cdot \right)$ *j k* $G_k(\mathbf{r}, t) = \sum_i Y_i g_k(\mathbf{k}, t; \rho, v)$ *g ,t; ρ,^υ* Ξ \equiv Ξ \sum \mathbf{r} , t) = $\sum Y_i \mathcal{Q}_L$ (K) r Use elaborate models (3d-layered soil half-space) to estimate the Green's function.

Models for multi-component earthquake ground motions Earthquake ground acceleration at any point can be resolved into three components along three orthogonal directions.

 $A(t) = iX_1(t) + jX_2(t) + kX_3(t)$ $\Psi(t) = i\theta_1(t) + j\theta_2(t) + k\theta_3(t)$ **TranslationRotation**

$$
X(t) = \begin{cases} X_1(t) \\ X_2(t) \\ X_3(t) \end{cases}; \theta = \begin{cases} \theta_1(t) \\ \theta_2(t) \\ \theta_3(t) \end{cases}
$$

$$
\Sigma(t) = \begin{cases} X(t) \\ \mathbf{F}(t) \end{cases}
$$
: Treat this as a vector random process
\n
$$
\langle \Sigma(t) \rangle = 0
$$

\nFocus attention on translations
\n
$$
C_{XX}(t) = \langle X^t(t) X(t) \rangle =
$$
\n
$$
\begin{bmatrix} \langle X_1^2(t) \rangle & \langle X_1(t) X_2(t) \rangle & \langle X_1(t) X_3(t) \rangle \\ \langle X_1(t) X_2(t) \rangle & \langle X_2^2(t) \rangle & \langle X_2(t) X_3(t) \rangle \\ \langle X_1(t) X_3(t) \rangle & \langle X_2(t) X_3(t) \rangle & \langle X_3^2(t) \rangle \end{bmatrix}
$$

 $X_i(t) = e_i(t) S_i(t); i = 1, 2, 3$

Consider
$$
X(t) = \begin{cases} X_1(t) \\ X_2(t) \\ X_3(t) \end{cases} = \begin{cases} e_1(t)S_1(t) \\ e_2(t)S_2(t) \\ e_3(t)S_3(t) \end{cases}
$$

\nwhere $S(t) = \begin{cases} S_1(t) \\ S_2(t) \\ S_3(t) \end{cases}$ is a stationary vector random process
\nwith zero mean and $e_1(t)$, $e_2(t)$ & $e_3(t)$ are deterministic
\nenvelope functions. We assume $e_i(t) = e(t)$; $i = 1, 2, 3$
\n $\langle X(t) X^t(t + \tau) \rangle = e(t) e(t + \tau) \langle S(t) S^t(t + \tau) \rangle$
\n $\Rightarrow R_{XX}(t, t + \tau) = e(t) e(t + \tau) R_{SS}(\tau)$
\n $\Rightarrow R_{XX}(t, t) = \underbrace{e^2(t)}_{\text{SUS}} R_{SS}(0) / \text{C}$
\nNote: $R_{SS}(0)$ is constant since $S(t)$ is stationary.
\nAlso, $R_{SS}(0)$ is symmetric and expected to be fully populated

$$
R_{XX}(t,t) = e^{2}(t) R_{SS}(0)
$$

We introduce a transformation

$$
\tilde{S}(t) = \Phi^{t} S(t)
$$

where Φ is a 3×3 transformation matrix.
Clearly, $\langle \tilde{S}(t) \rangle = 0$ and
 $\langle \tilde{S}(t) \tilde{S}^{t}(t + \tau) \rangle = \langle \Phi^{t} S(t) S^{t}(t) \Phi \rangle$
 $\Rightarrow R_{\tilde{S}\tilde{S}}(\tau) = \Phi^{t} R_{SS}(\tau) \Phi$
 $\Rightarrow R_{\tilde{S}\tilde{S}}(0) = \Phi^{t} R_{SS}(0) \Phi$
Select Φ such that $\Phi^{t} R_{SS}(0) \Phi$ is diagonal.
 $\Rightarrow R_{\tilde{S}\tilde{S}}(0) = Diag[R_{11} \quad R_{22} \quad R_{33}]$
 $\Rightarrow \Phi$: matrix of eigenvectors of $R_{SS}(0)$.
 $\Rightarrow \overline{X}(t) = e(t) \Phi^{t} S(t) \Rightarrow R_{\overline{X} \tilde{X}}(0) = e^{2}(t) R_{\tilde{S}\tilde{S}}(0)$

 $(t) = e_2(t) = e_3(t)$ $\big(t\big), \tilde{S}_{2}\big(t\big), \text{ and } \tilde{S}_{3}\big(t\big)$ The psd matirx of $\tilde{S}(t) = \left\{ \tilde{S}_1(t) \ \tilde{S}_2(t) \ \tilde{S}_3(t) \right\}^t$ is diagonal $_{\tilde{\rm SS}}\left(0\right)$ is diagonal does not imply that psd matirx $1\binom{v}{2}$ $\binom{v}{2}$ $\binom{v}{3}$ $1\binom{2}{3}$ $2\binom{2}{7}$, and $\frac{2}{3}$ (t) $S_2(t)$, and $S_3(t)$: are uncorrelated random processes \Rightarrow The psd matirx of $\tilde{S}(t) = \{\tilde{S}_1(t) \mid \tilde{S}_2(t) \mid \tilde{S}_3(t)\}$ $e_1(t) = e_2(t) = e_3(t) = e(t)$ $S_1(t), S_2(t), \text{ and } S_2(t)$ $R_{\tilde{g}\tilde{g}}(0)$ is diagonal does not imply that psd matirx of $\bullet e_1(t) = e_2(t) = e_2(t) =$ \bullet **Simplifying assumptions Note :** $\tilde{\zeta}(t)$ $\tilde{\zeta}(t)$ and $\tilde{\zeta}$ $\tilde{\mathcal{S}}(t) = \int \tilde{\mathcal{S}}(t) \tilde{\mathcal{S}}(t) \tilde{\mathcal{S}}(t)$ $\left(t\right)=\left\{\tilde{S}_{1}\left(t\right)\ \tilde{S}_{2}\left(t\right)\ \tilde{S}_{3}\left(t\right)\right\}$ $\big(t\big), \tilde{S}_{2}\big(t\big), \text{ and } \tilde{S}_{3}\big(t\big)$ $1²$) \sim 2 $²$) \sim 3</sup> $1\binom{2}{7}$, $2\binom{2}{7}$, and $2\binom{3}{4}$ $S_2(t) S_3(t)$ is diagonal. If $S_1(t)$, $S_2(t)$, and $S_3(t)$ are broadbanded, the above assumption can be deemed to be reasonable. $\tilde{S}(t) = \left\{ \tilde{S}_1(t) \right\}^t \tilde{S}_2(t) \left\{ \tilde{S}_2(t) \right\}^t$ $S_1(t), S_2(t), \text{ and } S_2(t)$ Ξ $\tilde{\mathbf{S}}(t) = \int \tilde{\mathbf{S}}(t) \tilde{\mathbf{S}}(t) \tilde{\mathbf{S}}(t)$ $\tilde{\zeta}(t)$ $\tilde{\zeta}(t)$ and $\tilde{\zeta}$

Principal axes of excitations

- Find the direction of coordinate axes in which the
- $R_{_{SS}}(0)$ matrix becomes diagonal.
- associated with the matrix $R_{\scriptscriptstyle{SS}}\left(0\right)$; This can be done by solving the eigenvalue problem *R*

• The major principal direction lies on the horizontal plane in a direction that points towards epicentre from the recording station. The minor direction is in the vertical plane.

This is an empirically observed feature from recorded data and there exists no "proof" for this.

Most often structures are designed taking into account only the horizontal components.

The principal directions for excitations need not coincide with the global coordinate axes used in modeling the structure.

For a structure that is symmetric in plan, excitation in one of the horizontal directions does not induce stresses in the other orthogonal direction.

Most structures are irregular in plan and the bendin g and torsional action could be coupled in the predominant modes of the struct ural oscillations. The modes could also be closely spaced.

• Under the action of earthquake ground motions the structures undergo significant torsional oscillations. This is one of the most characterist ic features of earthquake response of structures.

The structure translates and twists.

Principal axes for excitations exist and the ground motion components are uncorrelated along these axes.

Stationary random vibration analysis and basis for developing modal combination rule

Ref: W Smeby and A Der Kiureghian, 1985, EESD, $13, 1-12.$

 $M\ddot{U} + C\dot{U} + KU = -ML\ddot{X}$

 $U =$ vector of nodal displacements relative to the ground

$$
X = \left| \underbrace{X_1(t) \, X_2(t)}_{\text{Horizontal}} \, \underbrace{X_3(t)}_{\text{Vertical}} \right| = \text{translational components}
$$

of ground motion

$$
L = -\begin{bmatrix} L_1 & L_2 & L_3 \end{bmatrix}^t = \text{influence matrix.}
$$

 $v(t) = q'U(t)$ = response quantity of interest.

$$
M\ddot{U} + C\dot{U} + KU = -ML\ddot{X}
$$

\nLet $U = \Phi Y$ with
\n $K\Phi = \Lambda M\Phi$, $\Phi' M\Phi = I$, $\Phi' K\Phi = \Lambda$, $& \tilde{C} = \Phi' C\Phi$ diagonal
\n $\Rightarrow I\ddot{Y} + \tilde{C}\dot{Y} + \Lambda Y = -\Phi'ML\ddot{X}(t)$
\n $H(\omega) = \left[\Lambda - I\omega^2 + i\omega \tilde{C}\right]^{-1}$: diagonal matrix
\n $Y(\omega) = -H(\omega)\Phi'ML\ddot{X}(\omega)$
\n $G_{YY}(\omega) = H(\omega)\Phi'ML G_{XX}(\omega)L'M\Phi H^{*t}(\omega)$
\n $G_{UU}(\omega) = \Phi G_{YY}(\omega)\Phi'$
\n $v(t) = q'U$
\n $G_{yy}(\omega) = q' G_{UU}(\omega)q$
\n $G_{yy}(\omega) = q' \Phi H(\omega)\Phi'ML G_{XX}(\omega)L'M\Phi H^{*t}(\omega)\Phi'q$

Generic response quantity:
$$
\underline{v(t)} = \sum_{k=1}^{3} \sum_{i=1}^{N} \underbrace{V_i(t)}_{\text{Parification factor of the x direction component}}
$$

\n
$$
G_w(\omega) = \sum_{k=1}^{3} \sum_{l=1}^{3} \sum_{j=1}^{N} \sum_{j=1}^{N} \Psi_i^{(k)} \Psi_j^{(l)} H_i(\omega) H_j(-\omega) G_{\ddot{X}_k \ddot{X}_l}(\omega)
$$

\nLet
$$
Z(t) = \left[Z_1(t) \quad Z_2(t) \quad Z_3(t) \right]^t
$$
 be the ground motion components along the principal axes and let

\n
$$
X(t) = AZ(t)
$$

\n
$$
\Rightarrow G_{\ddot{X}\ddot{X}}(\omega) = AG_{\ddot{Z}\ddot{Z}}(\omega) A^t
$$

\nwhere $G_{\ddot{Z}_k \ddot{Z}_l}(\omega)$ is diagonal.

$$
X(t) = \begin{bmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix} Z(t) / / / X_1(t) = Z_1 \cos \theta + Z_2 \sin \theta \Rightarrow \langle X_1(t) \rangle = 0 X_2(t) = -Z_1 \sin \theta + Z_2 \cos \theta \Rightarrow \langle X_2(t) \rangle = 0 \sigma_1^2 = \text{Var} [X_1(t)] = \langle Z_1^2 \rangle \cos^2 \theta + \langle Z_2^2 \rangle \sin^2 \theta \sigma_2^2 = \text{Var} [X_2(t)] = -\langle Z_1^2 \rangle \sin^2 \theta + \langle Z_2^2 \rangle \cos^2 \theta \sigma_{12} = \langle X_1(t) X_2(t) \rangle = \langle Z_2^2 \rangle - \langle Z_1^2 \rangle \cos \theta \sin \theta \rho_{12} = \frac{\sigma_{12}}{\sqrt{\sigma_1^2 \sigma_2^2}} = \frac{-(1-\alpha)\sin 2\theta}{\sqrt{(1+\alpha)^2 - (1-\alpha)^2 \cos^2 2\theta}}; \alpha = \frac{\langle Z_2^2 \rangle}{\langle Z_1^2 \rangle} \end{bmatrix}
$$

 $\boldsymbol{f}(\boldsymbol{\omega})$ = $\boldsymbol{q}^t \boldsymbol{\Phi} H\big(\boldsymbol{\omega}\big) \boldsymbol{\Phi}^t MLA\big(\boldsymbol{\theta}\big) G_{\breve{\mathcal{I}} \breve{\mathcal{I}}} \big(\boldsymbol{\omega}\big) A^t \big(\boldsymbol{\theta}\big) L^t M\boldsymbol{\Phi} H^{*_t}\big(\boldsymbol{\omega}\big) \boldsymbol{\Phi}^t$ Spectral moments: $\lambda_{m} = \int \omega^{m} G_{vv} (\omega)$ Diagonal 0 $G_{\scriptscriptstyle\rm TV}\left(\omega\right)$ = $q^{\scriptscriptstyle I}\Phi H\left(\omega\right)\Phi^{\scriptscriptstyle I} MLA(\theta)G_{\scriptscriptstyle\!}Z\!\!\!Z\left(\omega\right)A^{\scriptscriptstyle I}\left(\theta\right)L^{\scriptscriptstyle I} M\Phi H^{{\scriptscriptstyle\rm T} \scriptscriptstyle I} \left(\omega\right)\Phi^{\scriptscriptstyle I} q$ \Rightarrow Spectral moments: $\lambda_m = \int \omega^m G_{vv}(\omega) d\omega$ ∞ \int $\overbrace{\hspace{4.5cm}}^{ }$ $\overline{}$

Leads to peak factors associated with mean and standard deviation of the maximum response over duration τ .

• One can determine the orientation θ for which the response variance reaches its maximum value.

• Alternatively, θ can be treated as a random variable and the expected values of response quantities of interest co uld be obtained with respect to pdf of θ .

Forms the basis for development of modal combination rule when the inputs are specfied in terms of a set of response spectra along the principal axes.

• When principal axes of excitation and structure axes coincide, or when excitation intensities along three axes are the same, a combination rule with

- SRSS for combination over excitation components, and
- \bullet CQC rule for combining over modal contributions can be obtained.

• More general forms which takes into account the value of θ have also been developed.

Earthquake source mechanism, wave propagation, site amplfication, and ground motion models

Earthquake ground motion=convolution of the source mechanism with the Green's funciton representing the wave pro pagation.

Application of double couples is equivalent to displacement discontinuities due to faulting.

 Boore and Atkinson, 1986, BSSA $\big(\omega \big)$ = $C S^{}_{1} \big(\omega \big) S^{}_{2} \big(\omega \big) S^{}_{3} \big(\omega \big)$ $S_1(\omega)$ = source spectrum $S_2(\omega)$ = amplification factor Fourier amplitude spectrum of ground acceleration ${\rm A}$ scaling factor $S_S(\omega) = CS_1(\omega)S_2(\omega)S_1$ *C* ω = CS₁ (ω | S₂ (ω | S₂ (ω) Ξ **PSD models based on seismological considerations** $S_3(\omega)$ = attenuation factor 3 1 4*r* $C = \frac{R_{\phi}FP_r}{4\pi\rho\beta^3} \left(\frac{1}{R}\right)$ ϕ $\pi\rho\beta$ (1) \equiv $=\frac{\sqrt[p]{r}}{4\pi\rho\beta^3}\left(\frac{R}{R}\right)$

$$
A_{s} (\omega) = CS_{1} (\omega) S_{2} (\omega) S_{3} (\omega)
$$

\n
$$
C = \frac{R_{\phi}FP_{r}}{4\pi\rho\beta^{3}} \left(\frac{1}{R}\right) /
$$

\nFactors
\n
$$
R_{\phi}: \text{ radiation pattern of the seismic wave}
$$

\n
$$
F: \text{ free surface effect}
$$

\n
$$
P_{r}: \text{ partition of energy into horizontal components}
$$

\n
$$
\rho: \text{ mass density}
$$

\n
$$
\beta: \text{ seismic wave velocity}
$$

\n
$$
R: \text{hyperbiral distance}
$$

$$
S_1(\omega) = m_0 \frac{\omega^2}{1 + \left(\frac{\omega}{\omega_c}\right)^2}
$$

\n
$$
m_0 = \text{seismic moment}
$$

\n
$$
\omega_c = \text{the frequency above which the spectral amplitudes of ground displacements begin to fall off; corner frequency (inversely proportional to source radius)
$$

\n
$$
f_e = \frac{2\pi}{\omega_c} = 0.49 \beta \left(\frac{\Delta \sigma}{m_0}\right)^{\frac{1}{3}}
$$

\n
$$
m_0 = 10^{(1.5M + 9.05)}/
$$

\n
$$
M = \text{earthquake moment magnitude}
$$

\n
$$
\Delta \sigma = \text{stress drop}
$$

\n(All quantities in SI units)

Kanai - Tajimi Model
\n
$$
S_2(\omega) = I \frac{(\omega_g^4 + 4\eta_g^2 \omega_g^2 \omega^2)}{(\omega^2 - \omega_g^2)^2 + 4\eta_g^2 \omega_g^2 \omega^2}
$$
\n
\nClough and Penzien model
\n
$$
S_2(\omega) = I \frac{(\omega_g^4 + 4\eta_g^2 \omega_g^2 \omega^2)}{(\omega^2 - \omega_g^2)^2 + 4\eta_g^2 \omega_g^2 \omega^2} \frac{H_f(\omega)^2}{High \text{ pass filter}}
$$
\n
$$
= I \frac{(\omega_g^4 + 4\eta_g^2 \omega_g^2 \omega^2)}{(\omega^2 - \omega_g^2)^2 + 4\eta_g^2 \omega_g^2 \omega^2} \frac{(\omega/\omega_f)^4}{[1 - (\omega/\omega_f)^2]^2 + 4\varsigma_f^2 (\omega/\omega_f)^2}
$$
\n
$$
\frac{High \text{ pass filter}}{High \text{ pass filter}}
$$

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 $J_3(\omega)$ = exp $-\frac{N\omega}{2\Omega\omega}$ $f(\omega)$ $f(\omega)$ = high cut filter dependent on f_m , a high cut-off frequency $\bigl(-\theta\omega\bigr)$ $(t) = e(t)F^{-1}(A_s(t))$ $\mathcal{L}(\omega,t) = e^2(t) \frac{1}{2 I} \left| A_s^2(\omega) \right| / 1 = \int |e(t)|^2$ 0 1 2 $\exp\left(-\frac{1}{2}\right)$ Q = quality factor of attenuation $= \exp$ (t) 1 $\left(\sqrt{\frac{2\pi}{T}} \right)^{1/s}$ *g* $\left(\begin{array}{ccc} 0 & 0 \\ 0 & 0 \end{array}\right)$ $\left(\begin{array}{ccc} -5 & 0 \\ -2 & 0 \end{array}\right)$ *gg* $\left(\begin{array}{c} \cdot & \cdot \\ \cdot & \cdot \end{array}\right)$ $\left(\begin{array}{c} \cdot & \cdot \\ \cdot & \cdot \end{array}\right)$ $\left(\begin{array}{c} \cdot & \cdot \\ \cdot & \cdot \end{array}\right)$ $S_3(\omega) = \exp\left(-\frac{R\omega}{2Q\beta}\right) f$ $X_{a}(t) = e(t)F^{-1}(A_{a}(t))$ $S_{gg}(\omega, t) = e^{2} (t) \frac{1}{2 \pi I} \left| A \right|$ $\omega(t) = e^{2}(t) - \frac{1}{2} \left[\omega \right] / 2 = |e(t)|^{2} dt$ ω) = exp $\left[-\frac{R\omega}{\omega}\right]$ $f(\omega)$ β π $\theta\omega$ ∞ $= e(t)F^{-}$ $(R\omega)$ $= \exp \left(-\frac{1}{2Q} \right)$ Ξ $\left(2Q\beta \right)$ Ξ $=\int$

Remark

The model relates the ground acceleration PSD to physical properties of the source and the medium through which the seismic waves travel.

Spatial variability of earthquake ground motions and response of multi - supported structures

•Long span bridges, large dams, pipelines, tunnels, Reference : A Zerva, 2009, Spatial variations of seismic gr ound motions, CRC Press, Boca Raton $\ddot{}$

SMART array at Taiwan

Two more stations at 2.8 and 4.8 km south of the centre.

Tri-axial accelerometers at every station

Why spatial variability occurs?

- Wave passage effect
- Extended source effect
- Scattering effect
- Attenuation effect

Questions

What are the phenomenological features associated with response of structures subjected to spatially varying ground motions?

• When it is important to consider them?

• How to model spatially varying ground motions as random processes?

• Based on data

• Based on phenomenological considerations How to develop modal combination rules when the inputs are specifed in terms of response sp ectra?