

# CRYOGENIC ENGINEERING



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Lecture No - **34**

## Earlier Lecture

- Cryogenic vessels use insulation to minimize all modes of heat transfer.
- Apparent thermal conductivity ( $k_A$ ) is calculated based on all possible modes of heat transfer.
- Expanded foam is a low density, cellular structure. A gas filled powder or a fibrous insulation reduces the gas convection due to the small size of voids.
- Radiation heat transfer is reduced by using radiation shields.

## Outline of the Lecture

### Topic : Cryogenic Insulation (contd)

- Vacuum
- Evacuated Powders
- Opacified Powders
- Tutorial

## Types of Insulation

- Expanded Foam – Mass
- Gas Filled Powders & Fibrous Materials – Mass
- Vacuum alone – Vacuum
- Evacuated Powders – Mass + Vacuum
- Opacified Powders – Mass + Vacuum + Reflective
- Multilayer Insulation – Vacuum + Reflective

## Introduction

- As seen earlier, the different modes of heat transfer are Conduction, Convection and Radiation.
- If the physical matter between the hot and the cold surfaces is removed, that is, by maintaining a perfect vacuum, Conduction and Convection are eliminated.
- However, Radiation heat transfer does not require any medium and in such cases, it is the only mode of heat transfer.

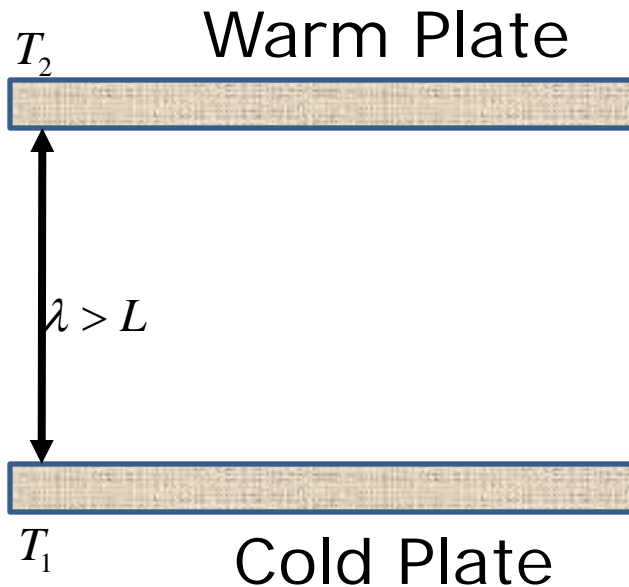
## Vacuum

- It is important to note that even in vacuum, there is some residual gas.
- These gas molecules contribute to the heat transfer by gaseous conduction.
- As the vacuum improves, this gas conduction decreases.
- In an ordinary conduction, a linear temperature gradient is built up. The molecules exchange heat with each other and as well as with the surfaces.

## Vacuum

- But in vacuum, the mean free path ( $\lambda$ ) of the molecules is more than the distance between the surfaces; the molecules rarely collide with each other.
- The energy is exchanged only between the surface and the colliding molecules.
- This type of heat transfer is called as free molecular conduction or residual gas conduction.
- This exists only at very low pressures or at very good vacuum.

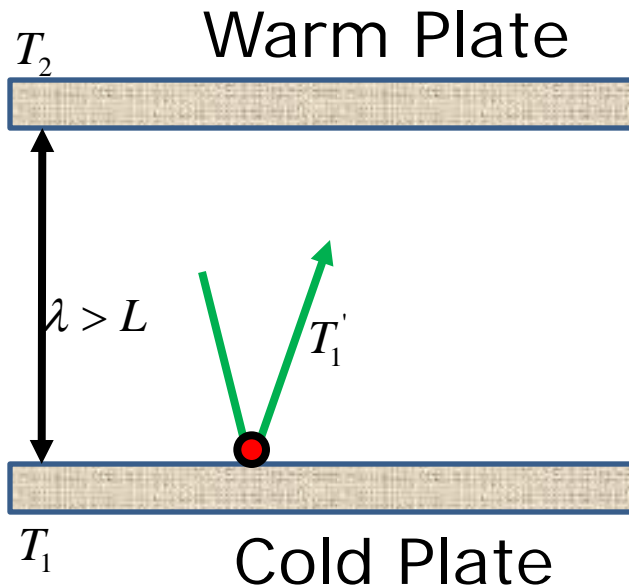
## Vacuum



- For the sake of understanding, consider two plates with temperatures  $T_1$  and  $T_2$ , ( $T_2 > T_1$ ) as shown.
- The gas pressure is very low in order to ensure that the mean free path ( $\lambda$ ) of the molecules is greater than  $L$ .
- In such situations, the gas molecules collide only with the surfaces and exchange energy.

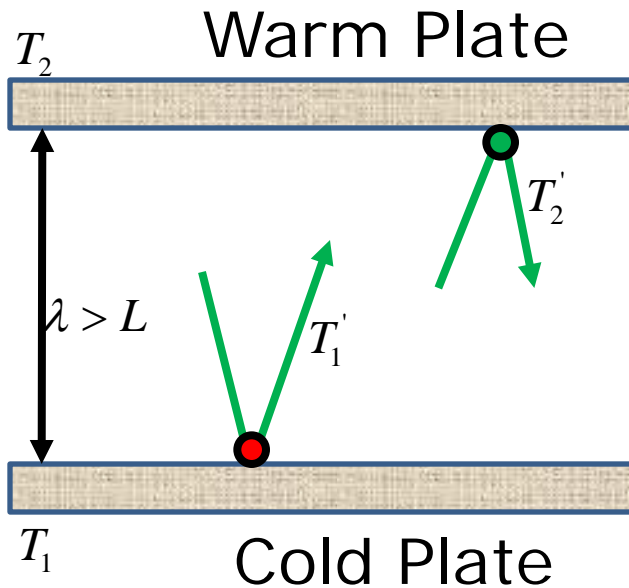


## Vacuum



- Consider a molecule colliding with bottom plate and leaving towards upper plate.
- The gas molecule collides with this surface at  $T_1$  and it transfers some energy to the surface.
- It leaves the cold surface with a kinetic energy corresponding to a temperature  $T'_1$ , higher than  $T_1$ .

## Vacuum

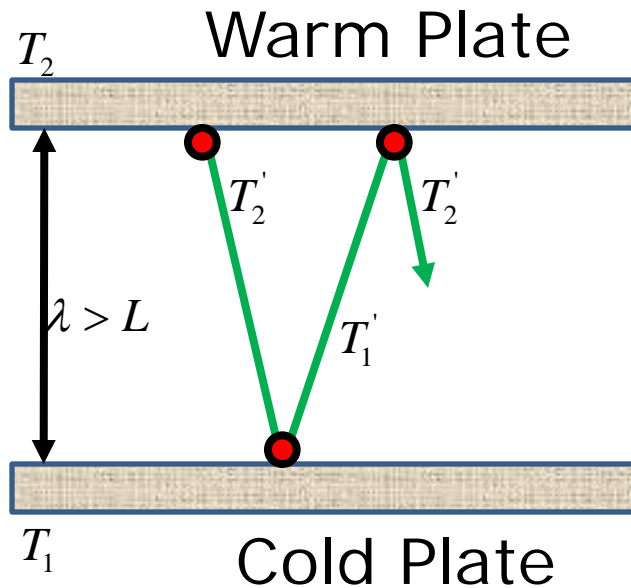


- Again, consider a molecule colliding with upper plate and leaving towards bottom plate.
- This gas molecule collides with surface at  $T_2$  and leaves at a temperature  $T_2'$ , lower than  $T_2$ .
- It is clear that, in both these impacts, thermal equilibrium is not attained. This process is repeated and contributes to free molecular conduction.

## Vacuum

- In order to measure the degree of thermal equilibrium between the molecule and the surface, we define Accommodation Coefficient (**a**).
- It is a ratio of actual energy transfer to the maximum possible energy transfer.
- Mathematically, 
$$a = \frac{\text{Actual Heat Transfer}}{\text{Max Heat Transfer}}$$
- Its value depends on the gas – surface interaction and the temperature of the surface.

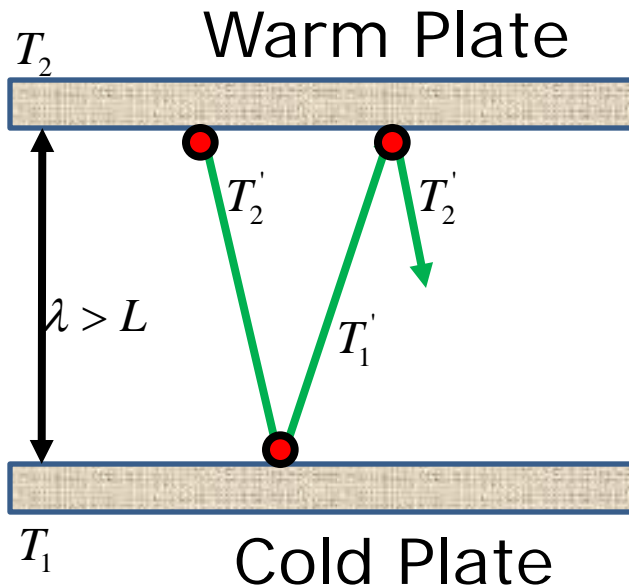
## Vacuum



- From the figure, for the cold surface, the actual temperature change is  $(T'_2 - T'_1)$ .
- But, the maximum possible temperature change is  $(T'_2 - T_1)$ .
- By definition, the accommodation coefficient for cold plate is

$$a_1 = \frac{T'_2 - T_1}{T'_2 - T'_1}$$

## Vacuum



- Similarly, for the hot surface, the actual temperature change is  $(T'_2 - T'_1)$ .
- But, the maximum possible temperature change is  $(T_2 - T'_1)$ .
- Therefore, the accommodation coefficient for the hot surface is given by

$$a_2 = \frac{T'_2 - T'_1}{T_2 - T'_1}$$

## Vacuum

- From the earlier slides, the accommodation coefficients are

$$a_1 = \frac{T_2' - T_1'}{T_2 - T_1}$$

$$a_2 = \frac{T_2' - T_1'}{T_2 - T_1}$$

- Rearranging the above equations, we have

$$T_1 = T_2' - \frac{T_2' - T_1'}{a_1}$$

$$T_2 = \frac{T_2' - T_1'}{a_2} + T_1'$$

$$T_2 - T_1 = (T_2' - T_1') \left( \frac{1}{a_1} + \frac{1}{a_2} - 1 \right)$$

## Vacuum

$$T_2 - T_1 = (T_2' - T_1') \left( \frac{1}{a_1} + \frac{1}{a_2} - 1 \right)$$

- Similar to an emissivity factor, we define a term accommodation factor  $F_a$ , which is given by

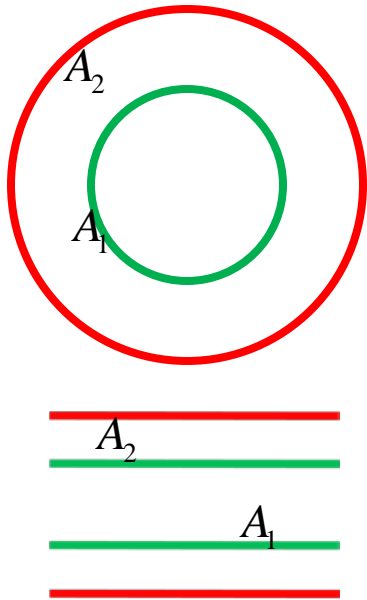
$$\frac{1}{F_a} = \left( \frac{1}{a_1} + \frac{1}{a_2} - 1 \right)$$

$$T_2 - T_1 = (T_2' - T_1') \frac{1}{F_a}$$

$$F_a = \frac{T_2' - T_1'}{T_2 - T_1}$$

## Vacuum

- The approximate accommodation coefficients for concentric sphere and concentric cylinder geometries are as tabulated below.



Temp (K)	He	H <sub>2</sub>	Ne	Air
300	0.29	0.29	0.66	0.8-0.9
78	0.42	0.53	0.83	1.0
20	0.59	0.97	1.0	1.0

- The subscript **1** denotes the enclosed surface and subscript **2** denotes the enclosure.



## Vacuum

Temp (K)	He	H <sub>2</sub>	Ne	Air
300	0.29	0.29	0.66	0.8-0.9
78	0.42	0.53	0.83	1.0
20	0.59	0.97	1.0	1.0

- At a given temperature, the accommodation coefficient increases with the increase in the molecular weight of the gas.
- For a given gas, the accommodation coefficient increases with the decrease in the temperature, due to better heat transfer at lower temperatures.

## Vacuum

- From the kinetic theory of gases, the total energy of a molecule is the sum of internal energy and kinetic energy.

- Mathematically,

$$e = U + KE$$

$$e = \left( c_v + \frac{R}{2} \right) T$$

- where,

$$\Delta e = \left( c_v + \frac{R}{2} \right) \Delta T$$

- R – Specific gas constant
- $c_v = R/(\gamma-1)$  – Specific heat of gas
- $\Delta T = (T'_2 - T'_1)$  – Change in temperature

## Vacuum

$$\Delta e = \left( c_v + \frac{R}{2} \right) \Delta T$$

- The definition of  $C_v$  and  $F_a$  are as given below.

$$c_v = \frac{R}{\gamma - 1}$$

$$T_2' - T_1' = F_a (T_2 - T_1)$$

- Substituting, we have

$$\Delta e = \left( \frac{R}{\gamma - 1} + \frac{R}{2} \right) (T_2 - T_1) F_a$$

$$\Delta e = \frac{F_a R}{2} (T_2 - T_1) \left( \frac{\gamma + 1}{\gamma - 1} \right)$$

## Vacuum

- The mass flux per unit time is given by

$$\frac{\dot{m}}{A} = \frac{\rho \bar{v}}{4}$$

- where,

- $\rho$  – Density,  $\bar{v}$  – Average velocity

- From Kinetic theory, average velocity is  $\bar{v} = \left( \frac{8RT}{\pi} \right)^{0.5}$

- Combining the above, together with equation of state, we have

$$\frac{\dot{m}}{A} = \frac{1}{4} \left( \frac{p}{RT} \right) \left( \frac{8RT}{\pi} \right)^{0.5}$$

$$\frac{\dot{m}}{A} = p \left( \frac{1}{2\pi RT} \right)^{0.5}$$

## Vacuum

- The total energy transfer per unit area owing to the molecular conduction is as given below.

$$\frac{\dot{Q}}{A} = \frac{\dot{m}}{A} (\Delta e) \quad \frac{\dot{m}}{A} = p \left( \frac{1}{2\pi RT} \right)^{0.5} \quad \Delta e = \frac{F_a R}{2} (T_2 - T_1) \left( \frac{\gamma + 1}{\gamma - 1} \right)$$

$$\frac{\dot{Q}}{A} = p \left( \frac{1}{2\pi RT} \right)^{0.5} \left( \frac{F_a R}{2} (T_2 - T_1) \left( \frac{\gamma + 1}{\gamma - 1} \right) \right)$$

$$\frac{\dot{Q}}{A} = \left( \left( \frac{\gamma + 1}{\gamma - 1} \right) \left( \frac{R}{8\pi T} \right)^{0.5} F_a \right) p (T_2 - T_1)$$

- T** is the temperature of the pressure gauge measuring the gas pressure.

## Vacuum

$$\frac{\dot{Q}}{A} = \left( \left( \frac{\gamma + 1}{\gamma - 1} \right) \left( \frac{R}{8\pi T} \right)^{0.5} F_a \right) p (T_2 - T_1)$$

- In the above equation, let us denote the term in the parenthesis by **G**. We have,

$$\dot{Q} = G p A (T_2 - T_1)$$

- Q** is valid only when the distance (**L**) between the plates is less than the mean free path (**λ**).

Mathematically,

$$L < \lambda = \frac{\mu}{p} \left( \frac{\pi RT}{2} \right)^{0.5}$$

## Vacuum

$$\dot{Q} = GpA(T_2 - T_1)$$

$$\lambda = \frac{\mu}{p} \left( \frac{\pi RT}{2} \right)^{0.5}$$

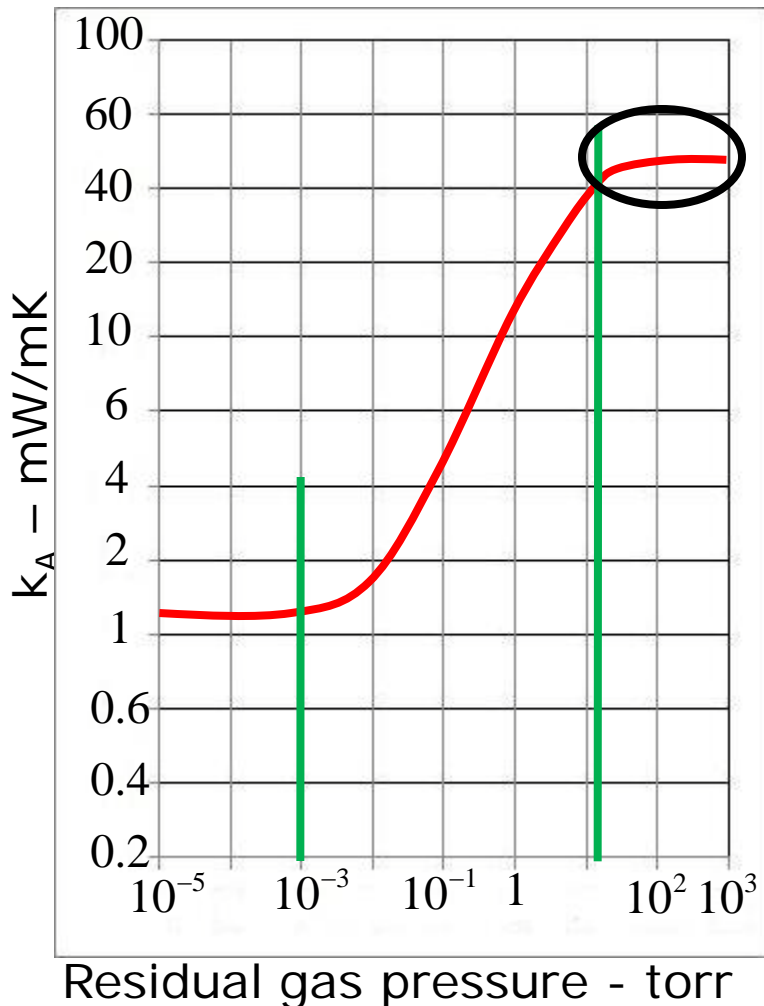
- From the above two equations, it is clear that the
  - The free molecular regime can be achieved by achieving very good vacuum.
  - The free molecular conduction heat transfer can be made negligible compared to other modes, by lowering the pressure, decreasing  $F_a$ , decreasing  $(T_2 - T_1)$ .

## Evacuated Powder

- Gas conduction is the primary and the dominant mode of heat transfer in a gas filled powder and fibrous insulations.
- One of the obvious ways to reduce this heat transfer is to evacuate the powder and the fibrous insulations.
- Usually, the vacuum that is commonly maintained in these insulations is in the range of  $10^3$  to  $10^{-5}$  torr. **1** torr = 1 mm of Hg.

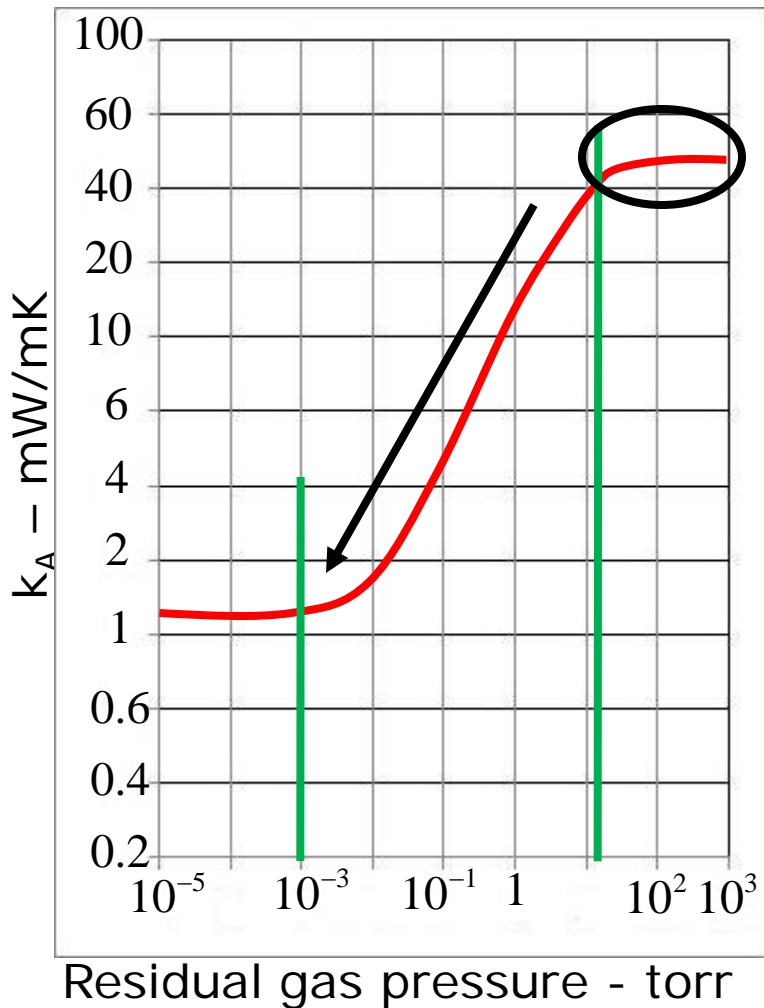


## Evacuated Powder



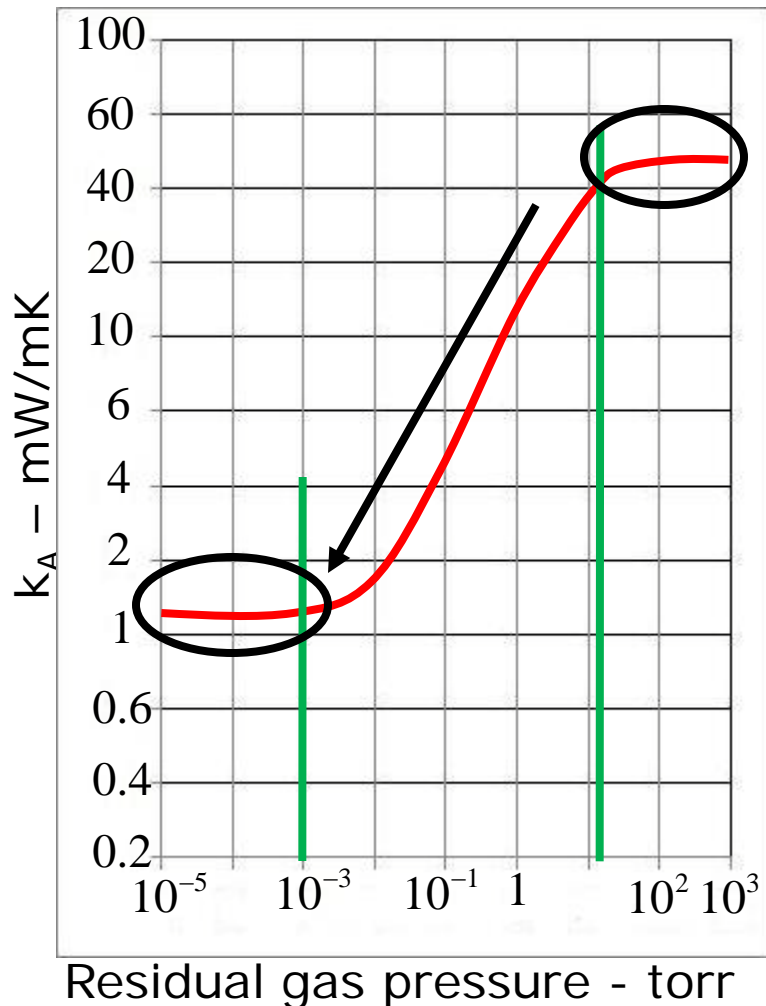
- The adjacent figure shows the variation of  $k_A$  with the residual gas pressure inside an evacuated powder insulation.
- $k_A$  is independent of residual gas pressures lying between atmospheric and 15 torr.

## Evacuated Powder



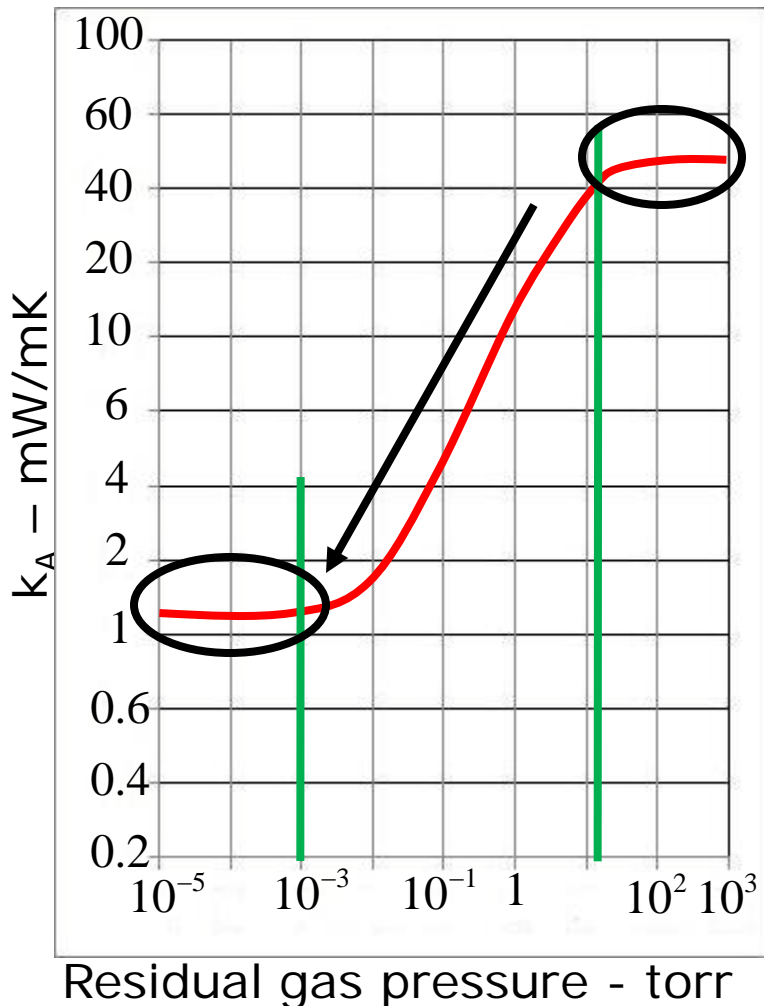
- With the lowering of pressure, 15 torr to  $10^{-3}$  torr,  $k_A$  becomes directly proportional to the pressure.
- It varies almost linearly on a logarithmic chart as shown.
- Here, the modes of heat transfer are due to radiation, solid conduction and free molecular conduction (dominant).

## Evacuated Powder



- With the further lowering of pressure, below  $10^{-3}$  torr, the variation of  $k_A$  is almost null.
- The mode of heat transfer is primarily due to solid conduction and radiation.
- Evacuated powders are superior in performance than vacuum alone in **300-77 K**, as the radiation heat transfer is comparatively less.

## Evacuated Powder



- At low pressures and temperatures, the solid conduction in evacuated powder dominates the radiant heat transfer.
- Hence, it is more advantageous to use vacuum alone in **77 K** to **4 K**. From Fourier's Law, we have

$$\dot{Q} = \frac{k_A A_m (T_h - T_c)}{\Delta x}$$

## Evacuated Powder

$$\dot{Q} = \frac{k_A A_m (T_h - T_c)}{\Delta x}$$

- where,
  - $k_A$  = Apparent thermal conductivity
  - $T_h - T_c$  = Temperature difference
  - $\Delta x$  = Distance
  - $A_m$  = Mean area of insulation.  $A_m$  for concentric cylinders and concentric spheres is as given below.

$$A_{m,cyl} = \frac{A_2 - A_1}{\ln \frac{A_2}{A_1}}$$

$$A_{m,sph} = \left( \frac{A_1 A_2}{A_1 + A_2} \right)^{\frac{1}{2}}$$

## Evacuated Powder

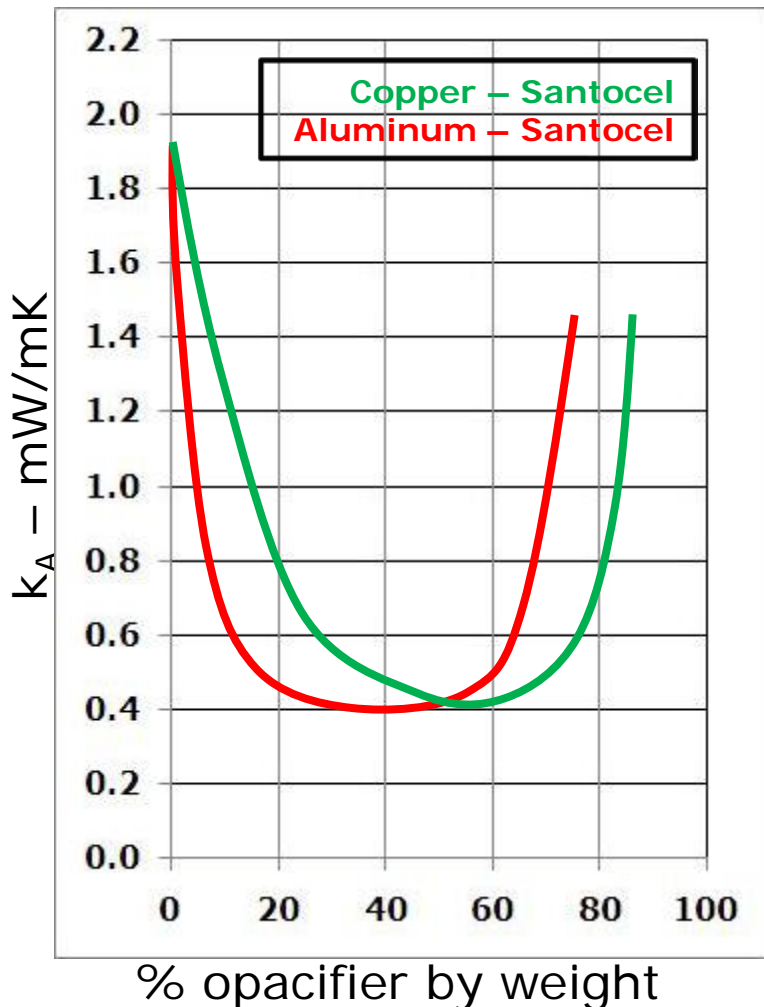
- The apparent thermal conductivity and density of few commonly used evacuated powder insulations are as shown.
- The residual gas pressure is less than  $10^{-3}$  torr for temperatures between **77 K** to **300 K**.

Powder	$\rho$ (kg/m <sup>3</sup> )	k (mW/mK)
Fine Perlite	180	0.95
Coarse Perlite	64	1.90
Lampblack	200	1.20
Fiberglass	50	1.70

## Opacified Powder Insulation

- Radiation heat transfer still contributes to the heat in leak in **300 K** to **77 K** temperature range in case of evacuated powders.
- In the year 1960, **Riede** and **Wang, Hunter et. al.** minimized this radiant heat transfer by addition of reflective flakes made of **Al** or **Cu** to the evacuated powder.
- These flakes act like radiant shields in the tiny heat transfer paths that are formed in the interstices of the evacuated powder.

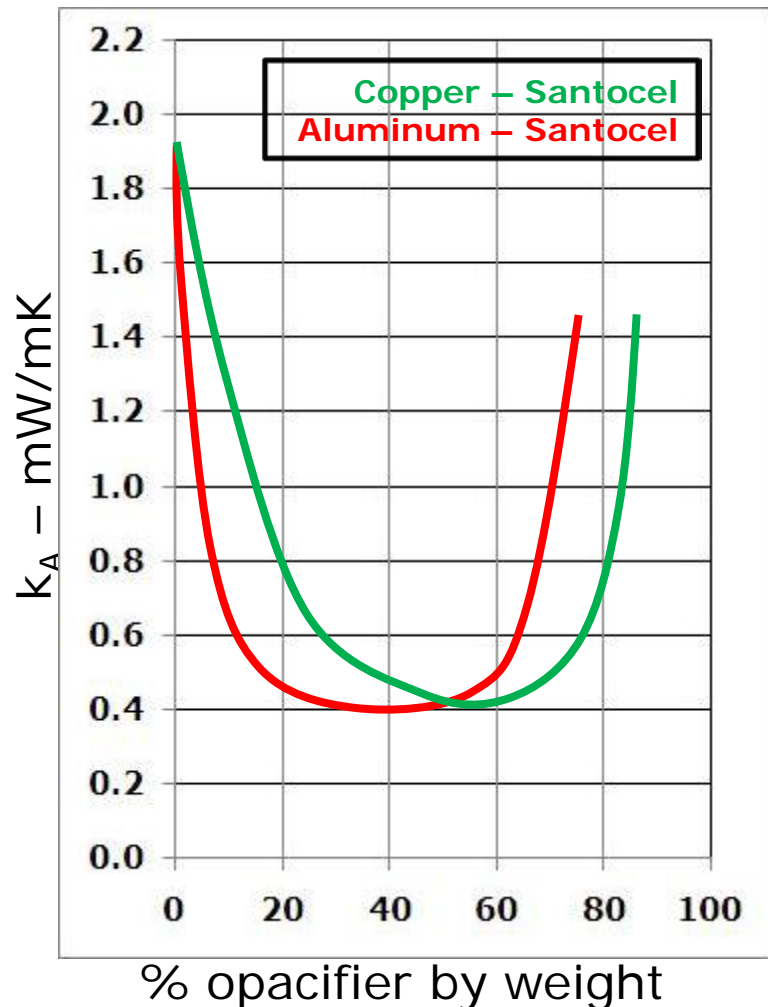
## Opacified Powder Insulation



- The figure shows the variation of % opacifier with thermal conductivity for **Cu – santocel** and **Al – santocel**.
- There exists an optimum operating point for each of these insulations.
- It has been observed that, with these additions,  $k_A$  can be reduced by 5 times.

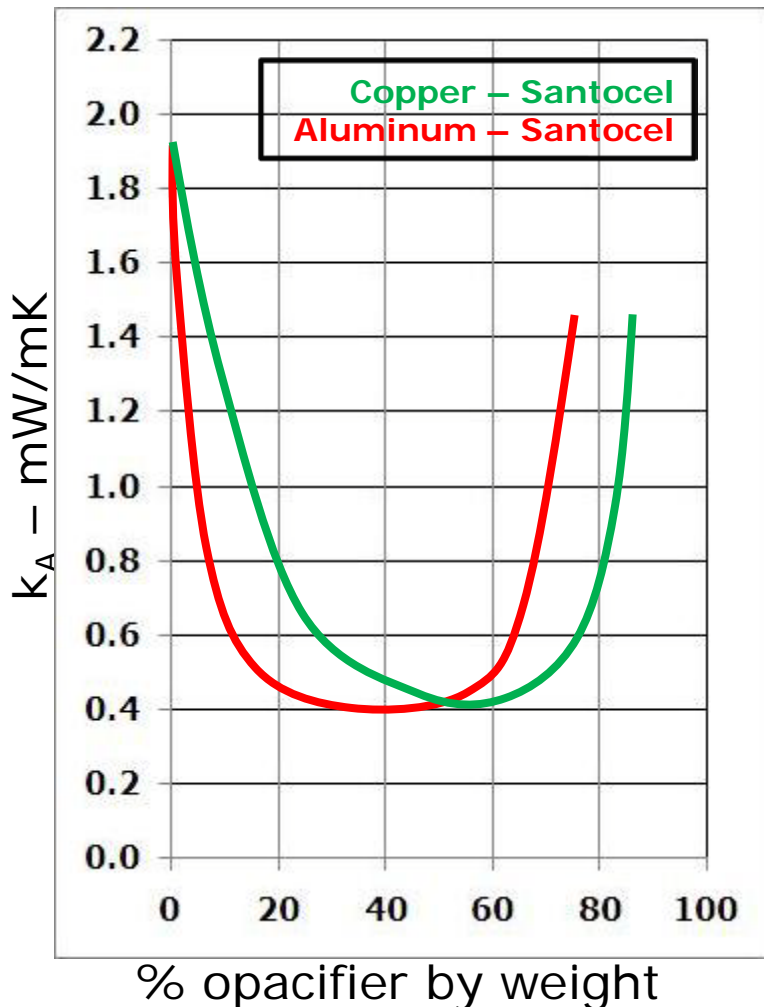


## Opacified Powder Insulation



- **Cu** flakes are more preferred as compared to **Al** flakes.
- The **Al** flakes have large heat of combustion.
- These together with **O<sub>2</sub>** can lead to accidents when used on **LOX** containers.

## Opacified Powder Insulation



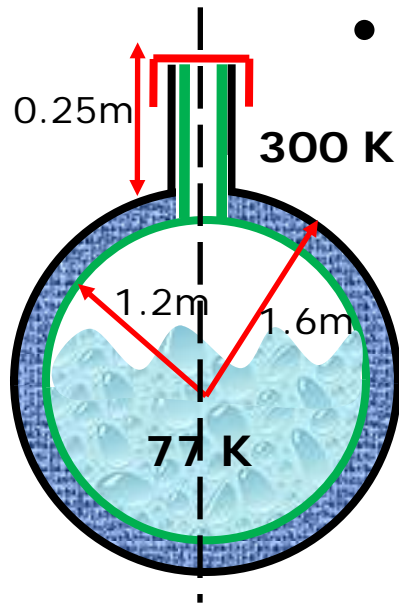
- Another disadvantage of this insulation is that the vibrations tend to pack the flakes together.
- This, not only increases the thermal conductivity but also short circuits the conduction heat transfer.

## Opacified Powder Insulation

- The apparent thermal conductivity (mW/mK) and density (kg/m<sup>3</sup>) of few commonly used opacified powder insulations are as shown.
- The residual gas pressure is less than 10<sup>-3</sup> torr for temperatures between **77 K** to **300 K**.

Powder	$\rho$ (kg/m <sup>3</sup> )	$k$ (mW/mK)
50/50 Cu – Santocel	180	0.33
40/60 Al – Santocel	160	0.35
50/50 Bronze – Santocel	179	0.58
Silica – Carbon	80	0.48

## Tutorial



- A spherical **LN2** vessel ( $\epsilon=0.8$ ) is as shown. The inner and outer radii are 1.2m and 1.6m respectively. Compare and comment on the heat in leak for the following cases.
  - Perlite (26 mW/mK)
  - Less Vacuum (1.5mPa)
  - Vacuum alone
  - Vacuum + 10 shields ( $\epsilon_s=0.05$ )
  - Evacuated Fine Perlite (0.95 mW/mK)
  - 50/50 Cu – Santocel (0.33 mW/mK)

## Tutorial

### Given

Apparatus : Spherical vessel ( $\epsilon=0.8$ )

Working Fluid : Liquid Nitrogen

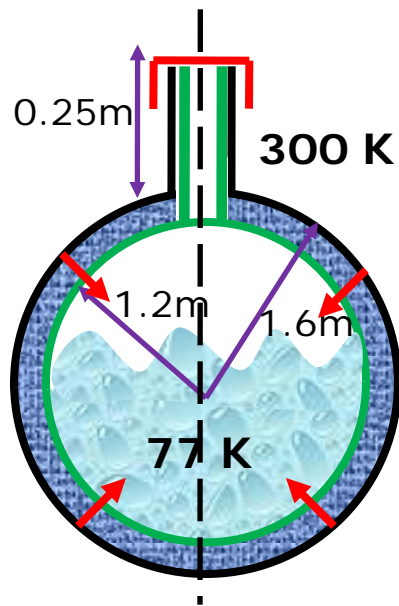
Temperature : 77 K (inner), 300 K (outer)

### Calculate heat in leak

- |   |                                     |
|---|-------------------------------------|
| 1 | Perlite (26 mW/mK)                  |
| 2 | Less Vacuum (1.5mPa)                |
| 3 | Vacuum alone                        |
| 4 | Vacuum + 10 shields                 |
| 5 | Evacuated Fine Perlite (0.95 mW/mK) |
| 6 | 50/50 Cu – Santocel (0.33 mW/mK)    |

- The shape factor between the two containers is assumed to be 1.

## Tutorial



Perlite ( $k_A = 26\text{mW/m-K}$ )

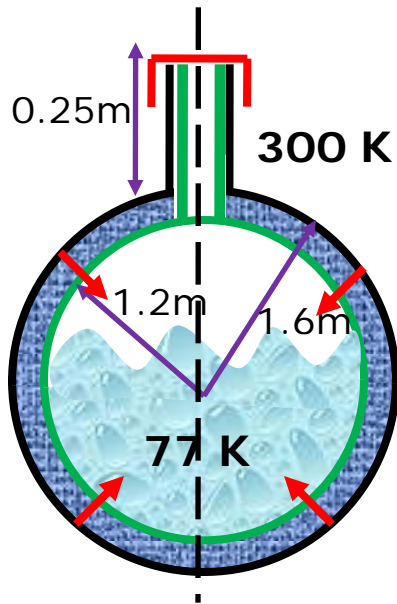
- **Sphere** -  $R_1 = 1.6\text{m}$ ,  $R_2 = 1.2\text{m}$ ,  $k_A$ ,  $\Delta T = (300 - 77) = 223$ .

$$Q = \frac{4\pi k_A R_1 R_2 \Delta T}{(R_2 - R_1)}$$

$$Q = \frac{4\pi (26)(10^{-3})(1.6)(1.2)(223)}{(1.6 - 1.2)}$$

$$Q = 349.7\text{W}$$

## Tutorial



Less Vacuum (1.5mPa)

- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  
 $e_1=e_2=0.8$ ,  $T_1=77\text{ K}$ ,  $T_2=300\text{ K}$ .

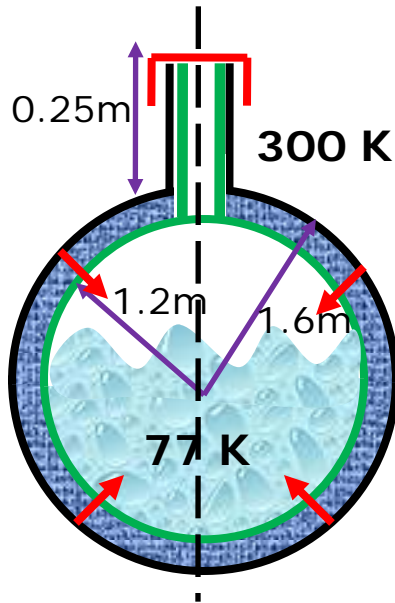
- The net heat transfer is due to both radiation and residual gas conduction.

$$F_e = \left( \frac{1}{e_1} + \left( \frac{A_1}{A_2} \right) \left( \frac{1}{e_2} - 1 \right) \right)^{-1}$$

$$F_e = \left( \frac{1}{0.8} + \left( \frac{1.2}{1.6} \right)^2 \left( \frac{1}{0.8} - 1 \right) \right)^{-1} = 0.72$$

## Tutorial

Less Vacuum (1.5mPa)



- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  
 $e_1=e_2=0.8$ ,  $T_1=77\text{ K}$ ,  $T_2=300\text{ K}$ .

$$Q = F_e F_{1 \rightarrow 2} \sigma A_1 (T_2^4 - T_1^4)$$

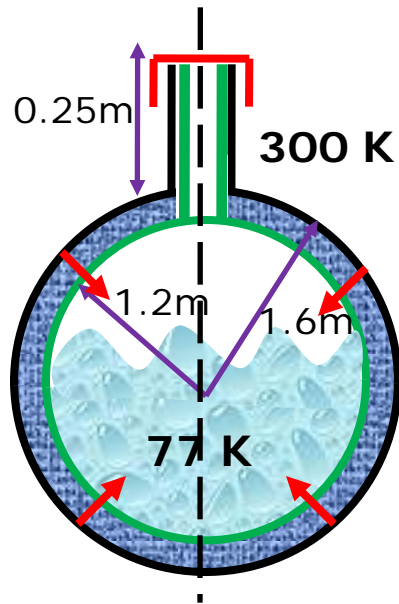
$$F_e = 0.72$$

$$Q = (0.72)(1)(5.67)(10^{-8})\pi(1.6^2)(300^4 - 77^4)$$

$$Q_r = 2648\text{W}$$



## Tutorial



Less Vacuum (1.5mPa)

- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  $T_1=77\text{ K}$ ,  $T_2=300\text{ K}$ ,  $p=1.5\text{ mPa}$ .

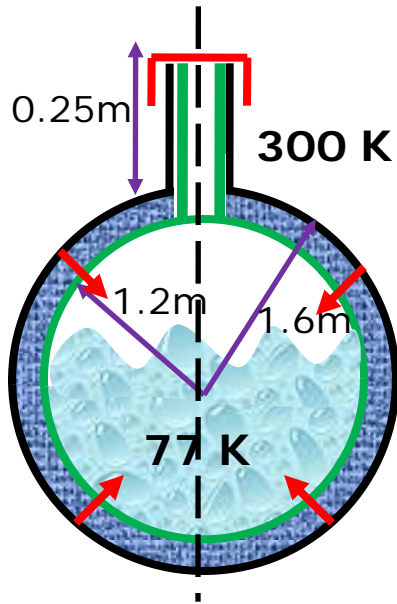
$$\lambda = \frac{\mu}{p} \left( \frac{\pi RT}{2} \right)^{0.5}$$

$$\lambda = \frac{(18.47)(10^{-6})}{(1.5)(10^{-3})} \left( \frac{\pi(287.6)(300)}{2} \right)^{0.5} = 4.53$$

- It is clear that the mean free path ( $\lambda$ ) is greater than distance between the surfaces (0.4m).

## Tutorial

Less Vacuum (1.5mPa)



- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  $T_1=77\text{ K}$ ,  $T_2=300\text{ K}$ ,  $p=1.5\text{ mPa}$ .

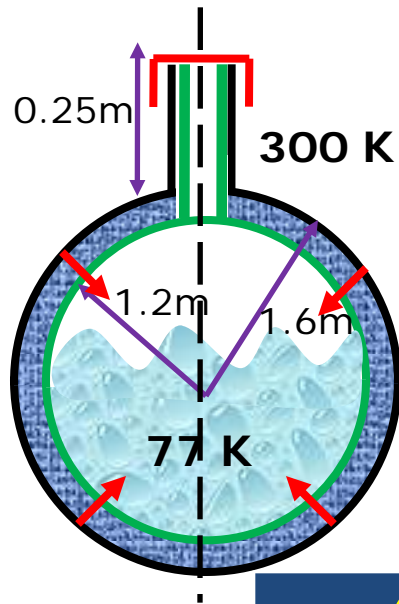
$$F_a = \left( \frac{1}{\alpha_1} + \left( \frac{A_1}{A_2} \right) \left( \frac{1}{\alpha_2} - 1 \right) \right)^{-1}$$

$$F_e = \left( \frac{1}{1} + \left( \frac{1.2}{1.6} \right)^2 \left( \frac{1}{0.85} - 1 \right) \right)^{-1} = 0.91$$

T (K)	Air
300	0.8-0.9
78	1.0
20	1.0

## Tutorial

Less Vacuum (1.5mPa)



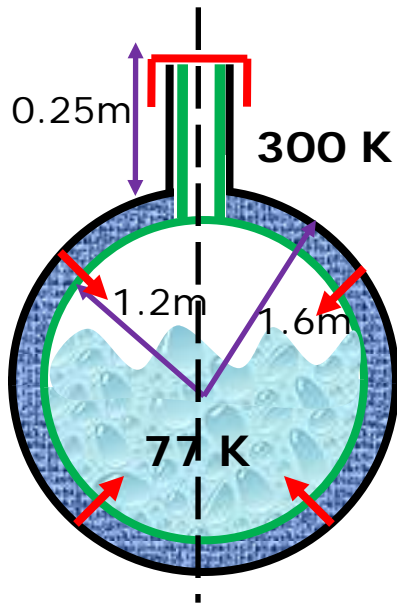
- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  $T_1=77\text{ K}$ ,  $T_2=300\text{ K}$ ,  $p=1.5\text{ mPa}$ .

$$\dot{Q} = \left( \left( \frac{\gamma + 1}{\gamma - 1} \right) \left( \frac{R}{8\pi T} \right)^{0.5} F_a \right) pA(T_2 - T_1)$$

$$\dot{Q} = \left( \left( \frac{1.4 + 1}{1.4 - 1} \right) \left( \frac{287.6}{8\pi (300)} \right)^{0.5} (0.91) \right) (1.5)(10^{-3})(300 - 77)$$

$$Q_{gc} = 0.356\text{W}$$

## Tutorial



Vacuum alone

- **Sphere** -  $R_1=1.6\text{m}$ ,  $R_2=1.2\text{m}$ ,  $k_A$ ,  $T_1=77\text{K}$ ,  $T_2=300\text{K}$ ,  $e_1$ ,  $e_2=0.8$ ,  $F_{1\rightarrow 2}=1$ .

$$Q = F_e F_{1\rightarrow 2} \sigma A_1 (T_2^4 - T_1^4)$$

$$F_e = 0.72$$

$$Q = (0.667)(1)(5.67)(10^{-8})\pi(1.6^2)(300^4 - 77^4)$$

$$Q = 2648\text{W}$$

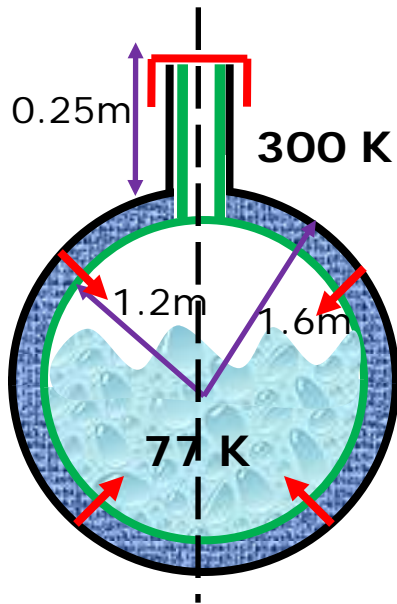
Vacuum + 10 shields

- $e_1$ ,  $e_2=0.8$ ,  $e_s=0.05$ .

$$F_e = 0.003$$

$$Q = 11.02\text{W}$$

## Tutorial



Evacuated Fine Perlite ( $k_A = 0.95 \text{ mW/mK}$ )

- **Sphere** -  $R_1 = 1.6 \text{ m}$ ,  $R_2 = 1.2 \text{ m}$ ,  $k_A$ ,  $\Delta T = (300 - 77) = 223$ .

$$Q = \frac{4\pi k_A R_1 R_2 \Delta T}{(R_2 - R_1)}$$

$$Q = 12.7 \text{ W}$$

50/50 Cu – Santocel ( $k_A = 0.33 \text{ mW/m-K}$ )

- **Sphere** -  $R_1 = 1.6 \text{ m}$ ,  $R_2 = 1.2 \text{ m}$ ,  $k_A$ ,  $\Delta T = (300 - 77) = 223$ .

$$Q = \frac{4\pi k_A R_1 R_2 \Delta T}{(R_2 - R_1)}$$

$$Q = 4.41 \text{ W}$$

## Tutorial

### Heat in leak (Q)

Perlite	349.7 W
Less Vacuum (1.5mPa)	$Q_r = 2648$ W $Q_{gc} = 0.356$ W
Vacuum alone	2648 W
Vacuum + 10 shields	11.02 W
Evacuated Fine Perlite	12.7 W
50/50 Cu – Santocel	4.41 W

## Summary

- In vacuum, the radiation is the dominant mode of heat transfer.
- Evacuated powders are superior in performance than vacuum alone in **300-77 K**, as the radiation heat transfer is comparatively less.
- At low pressures and temperatures, the solid conduction in evacuated powder dominates the radiant heat transfer.
- In an opacified powder, the radiation heat transfer is minimized by addition of reflective flakes.

**Thank You!**